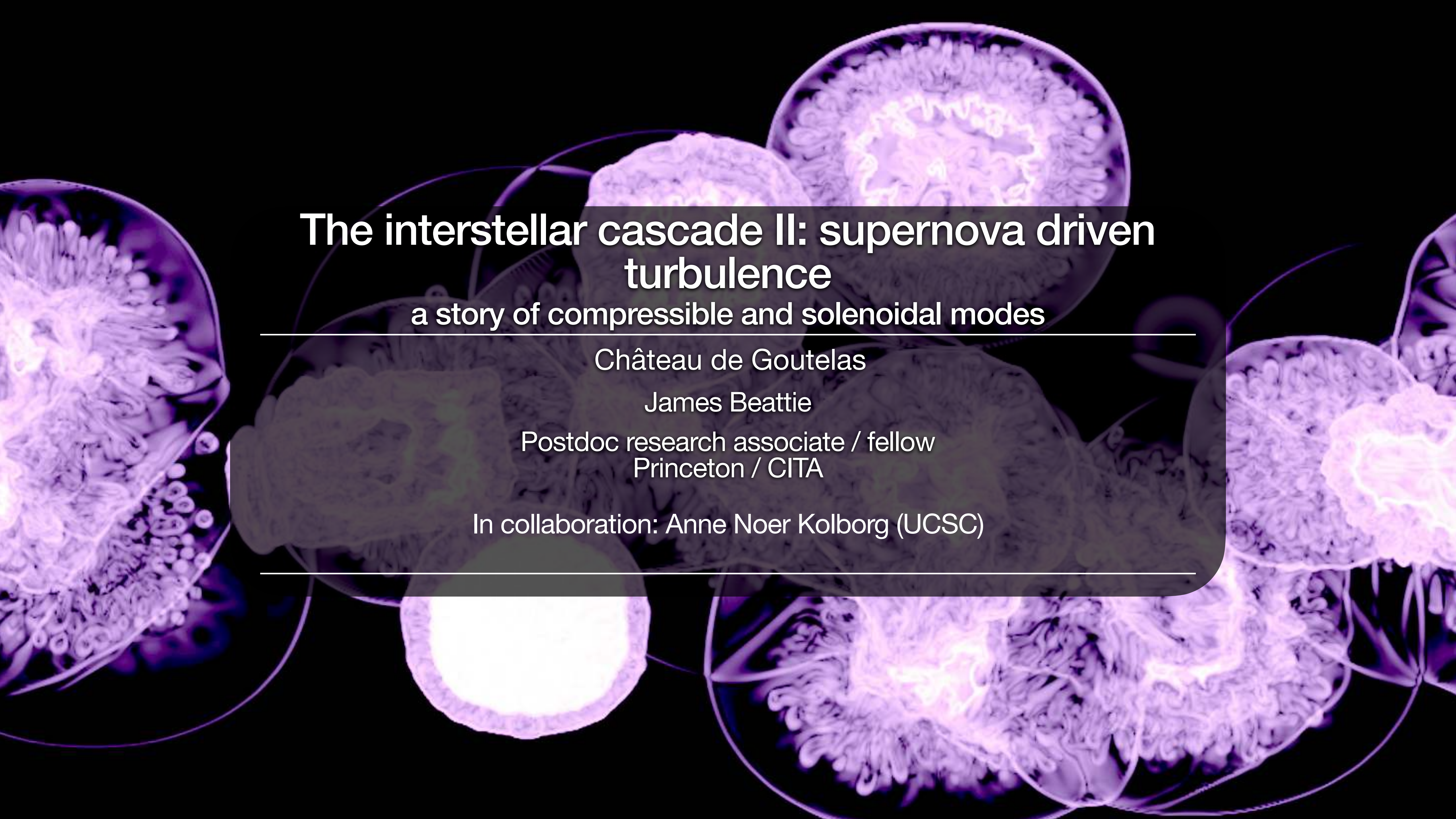




The interstellar cascade II: supernova driven turbulence

a story of compressible and solenoidal modes

Château de Goutelas

James Beattie

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Princeton / CITA

In collaboration: Anne Noer Kolborg (UCSC)

A beautiful degeneracy... or universality?

Is The Starry Night Turbulent?

JAMES R. BEATTIE¹ AND NECO KRIEL²

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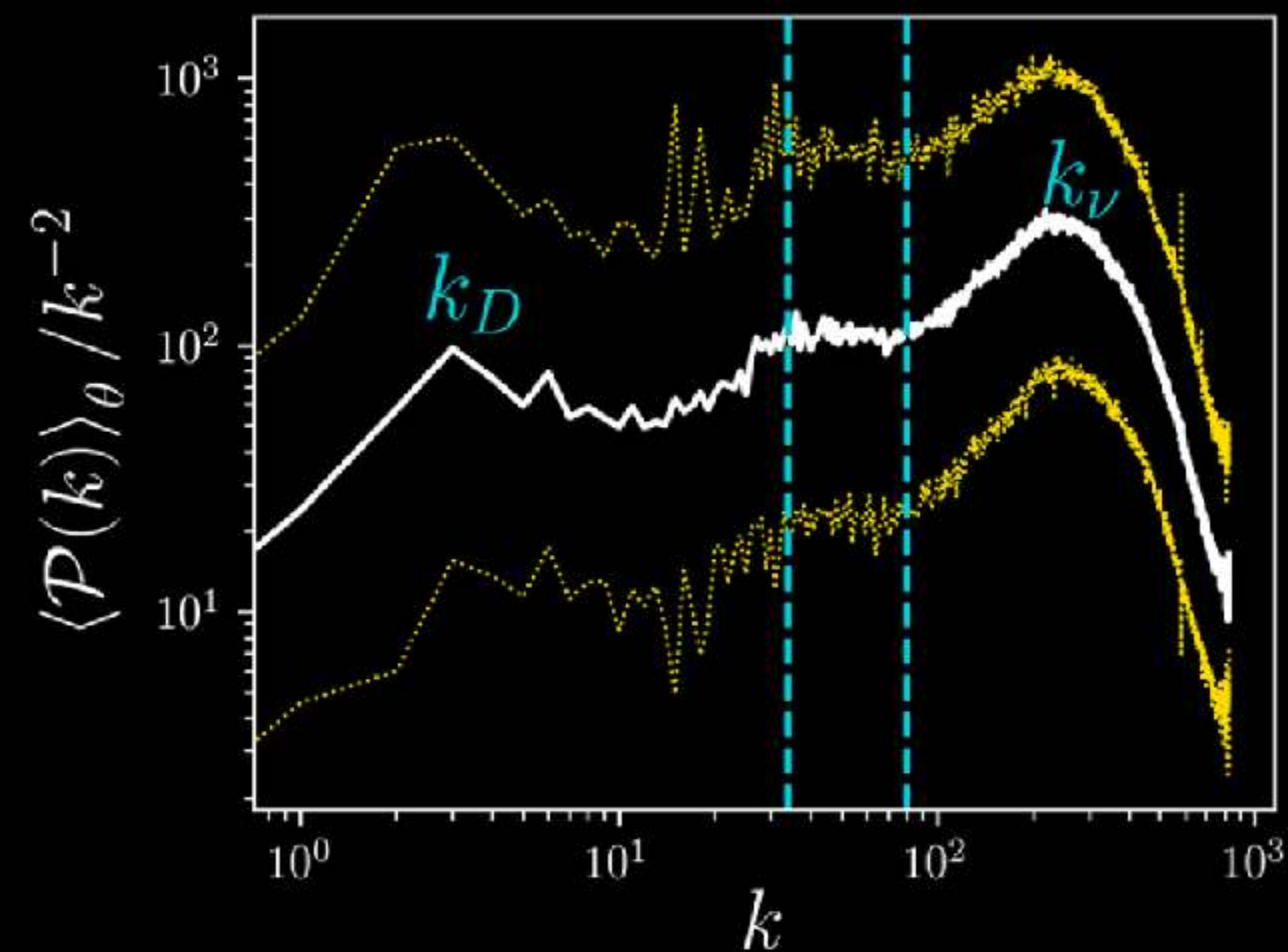
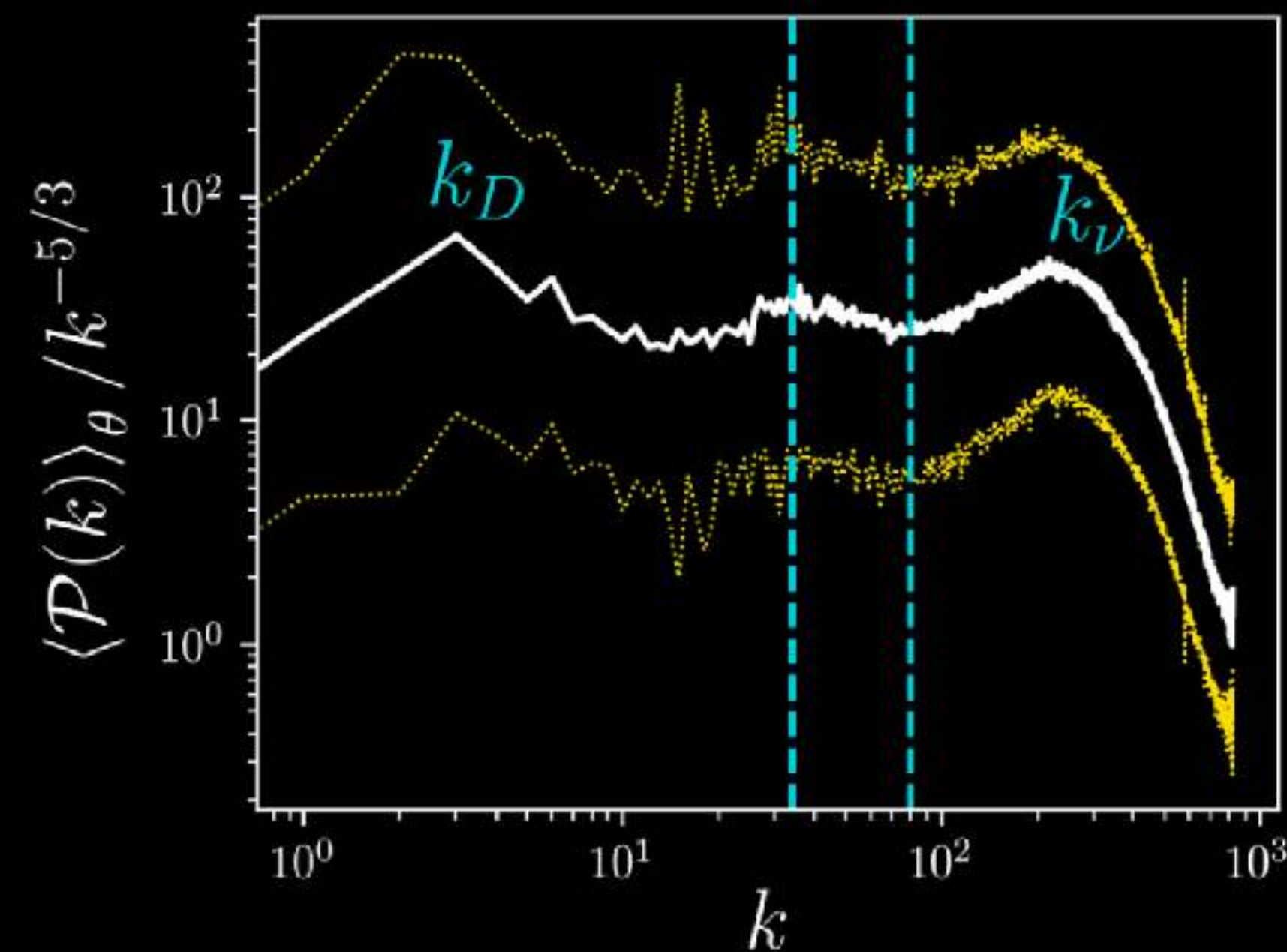
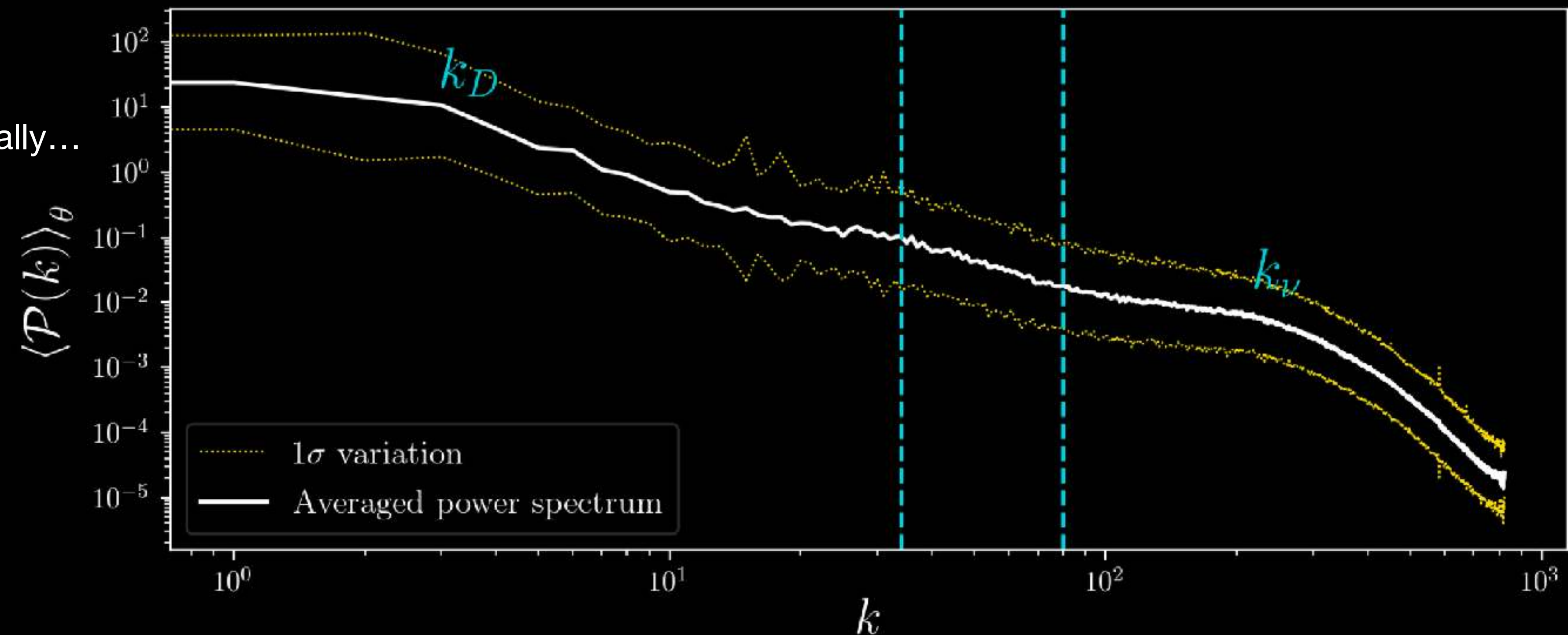
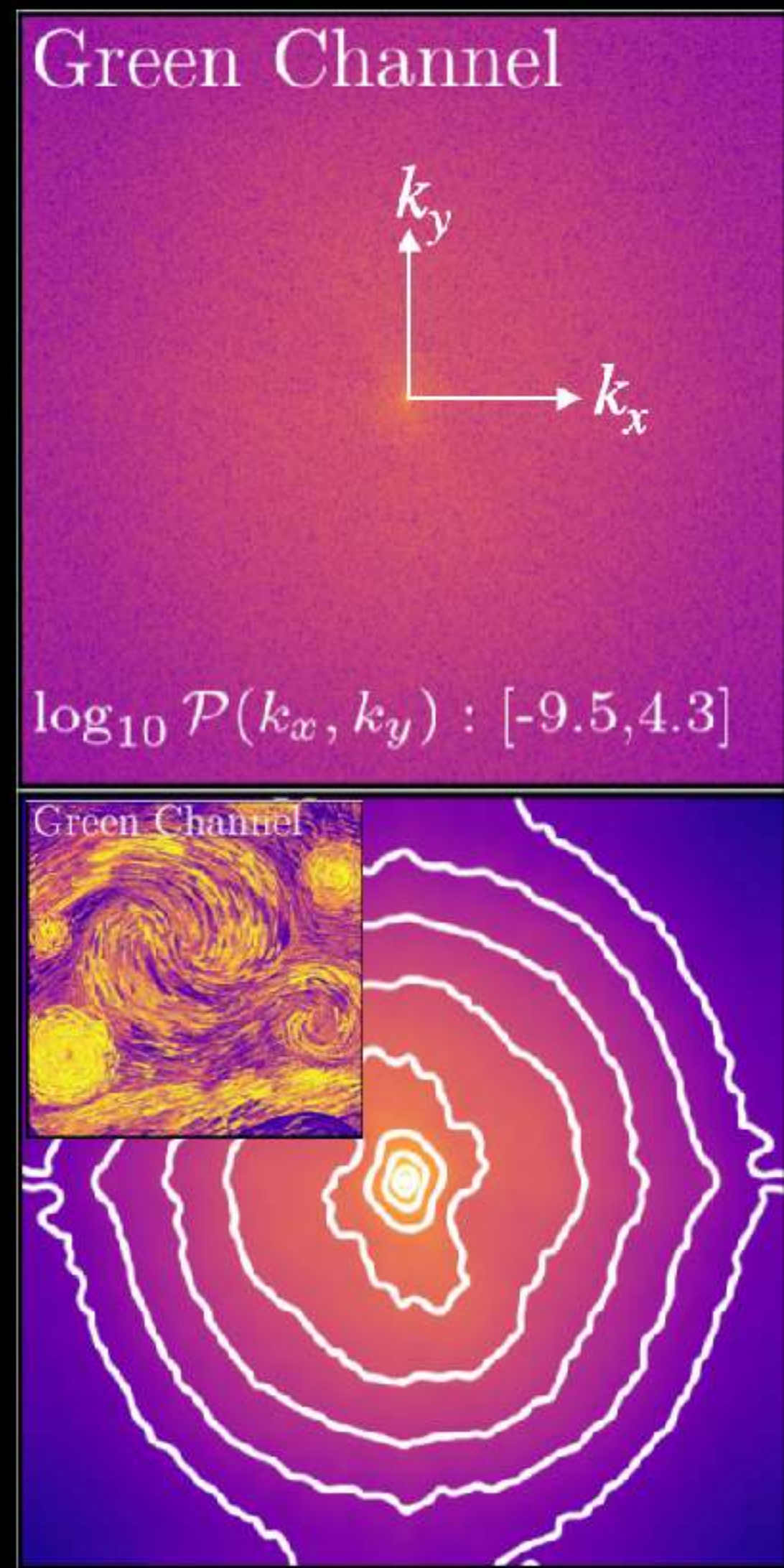
²*Science and Engineering Faculty, Queensland University of Technology, Brisbane, Australia*

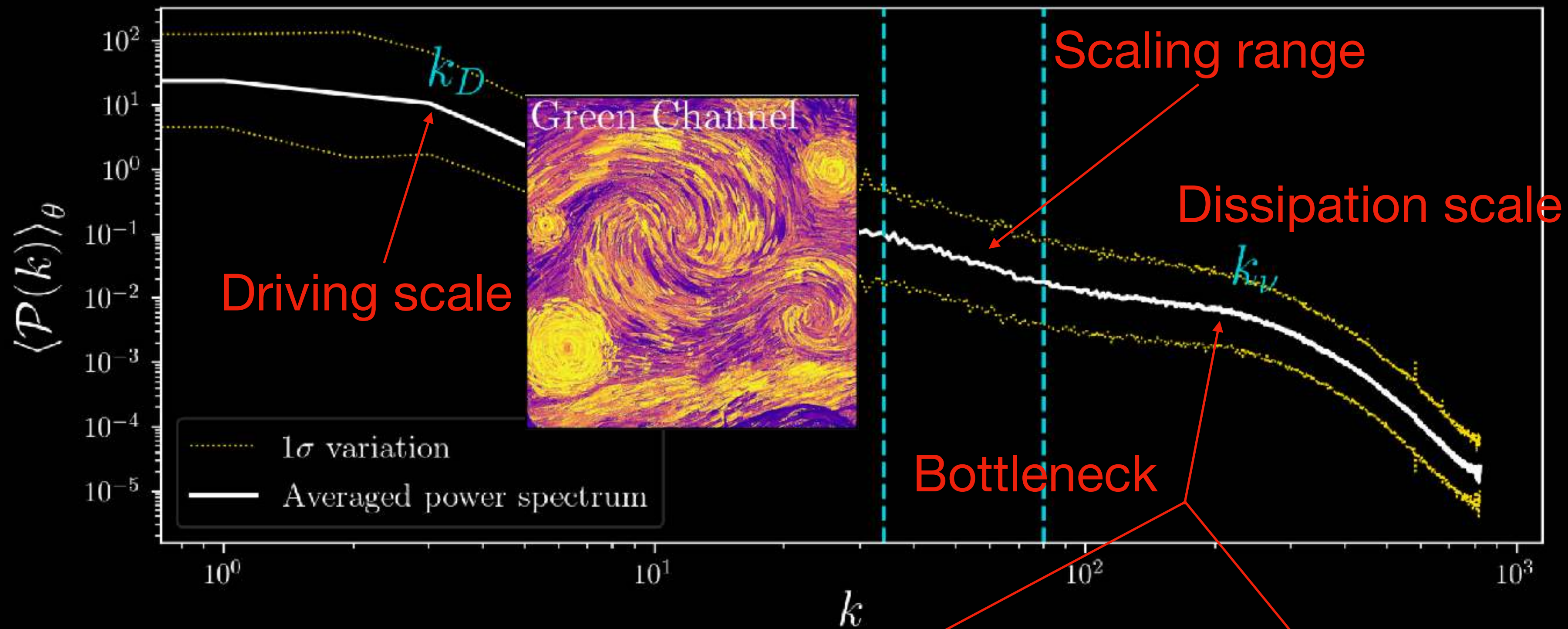
Vincent van Gogh's painting, *The Starry Night*, is an iconic piece of art and cultural history. The painting portrays a night sky full of stars, with eddies (spirals) both large and small. **Kolmogorov (1941)**'s description of subsonic, incompressible turbulence gives a model for turbulence that involves eddies interacting on many length scales, and so the question has been asked: is *The Starry Night* turbulent? To answer this question, we calculate the azimuthally averaged power spectrum of a square region (1165×1165 pixels) of night sky in *The Starry Night*. We find a power spectrum, $\mathcal{P}(k)$, where k is the wavevector, that shares the same features as supersonic turbulence. It has a power-law $\mathcal{P}(k) \propto k^{-2.1 \pm 0.3}$ in the scaling range, $34 \leq k \leq 80$. We identify a driving scale, $k_D = 3$, dissipation scale, $k_\nu = 220$ and a bottleneck. This leads us to believe that van Gogh's depiction of the starry night closely resembles the turbulence found in real molecular clouds, the birthplace of stars in the Universe.



back in 2017...

Van Gogh painted quite isotropically...

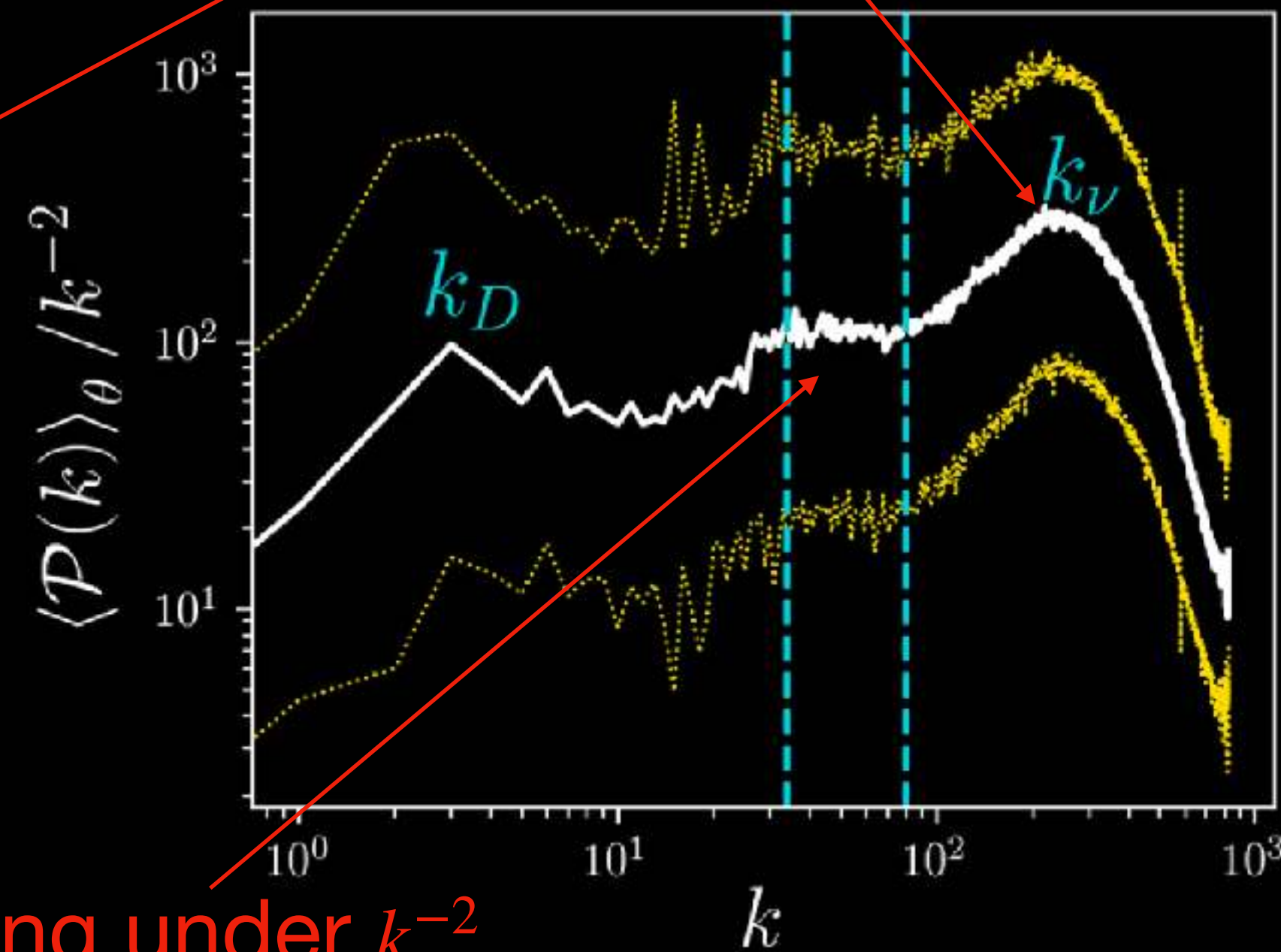
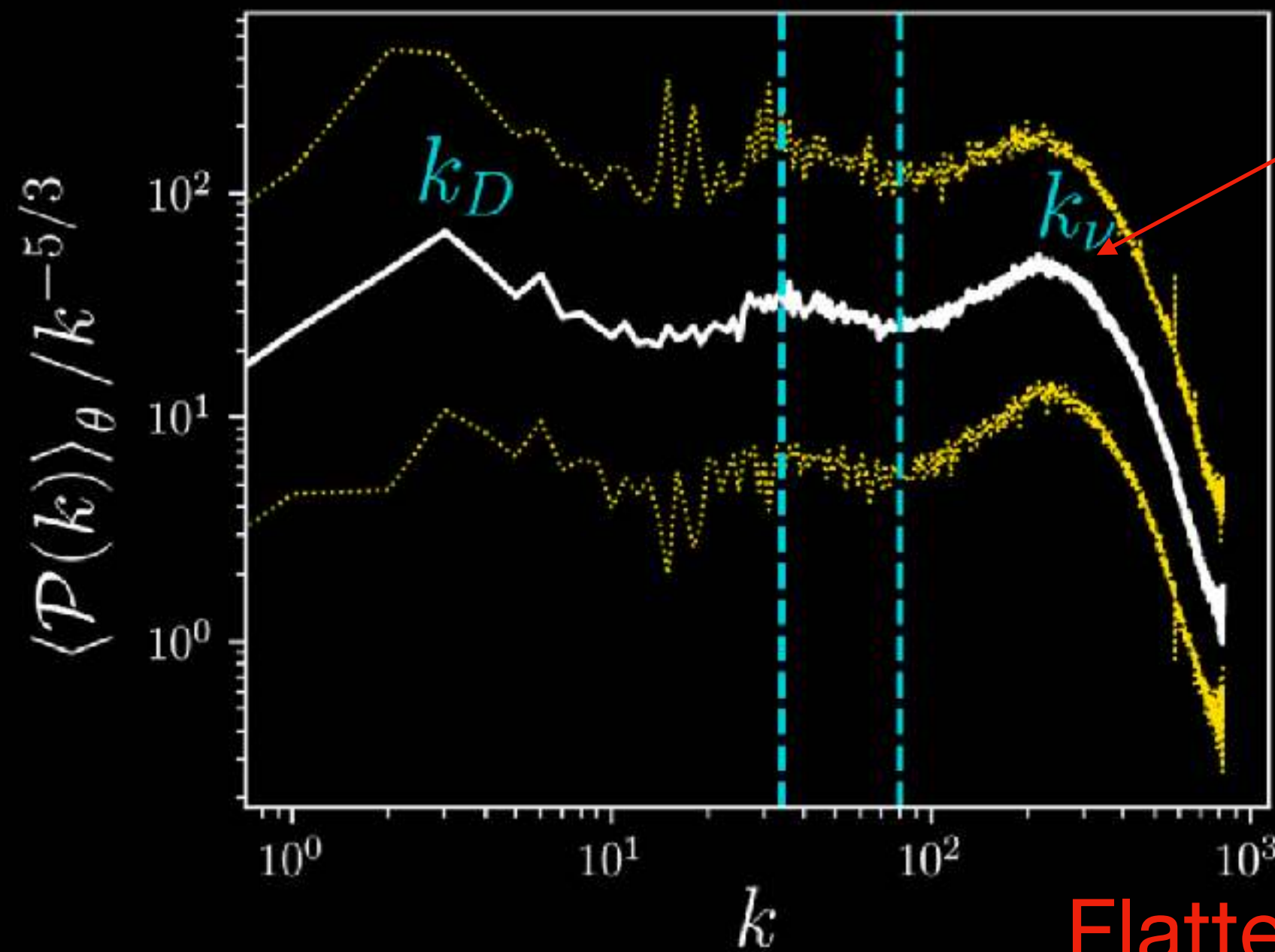




A beautiful degeneracy...
or universality?

$$k_\nu \approx 220 \text{ px}^{-1} \approx \frac{2\pi}{\ell_\nu}$$

$$\ell_\nu \approx (1 - 2) \text{ cm}$$



Flattening under k^{-2}

Objectives

1. Understand the nature of supernova-driven turbulence to directly see how much it resembles our simple Kolmogorov-style models (*philosophy: build the simplest models first, understand them in great detail*).
2. Introduce the idea of energy flux density statistics, defining some new opportunities for using these directly on observations

My four horsemen of ISM physics

Turbulent
magnetic fields

GeV – TeV
Cosmic rays

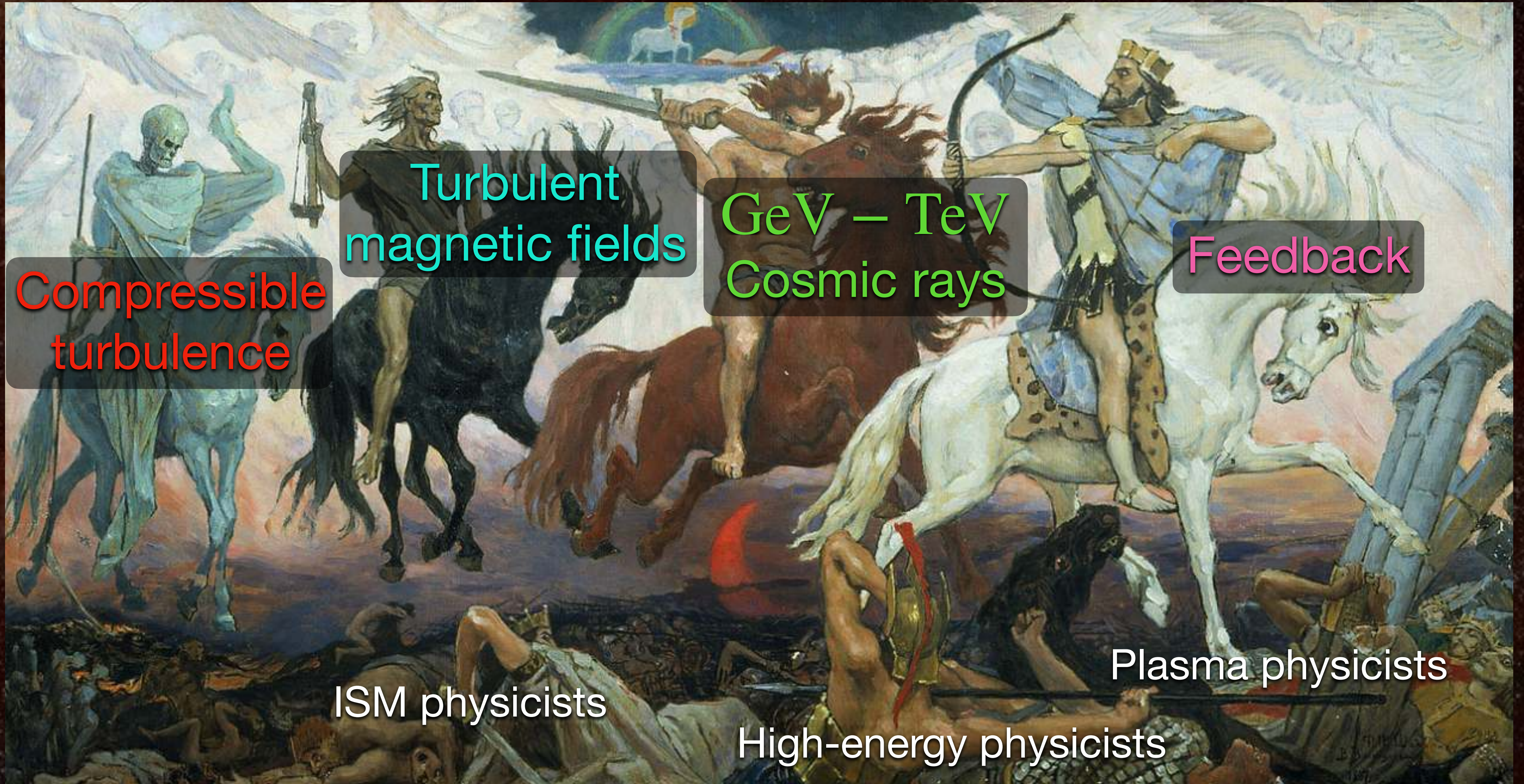
Feedback

Compressible
turbulence

ISM physicists

Plasma physicists

High-energy physicists



My four horsemen of ISM physics

Compressible
turbulence

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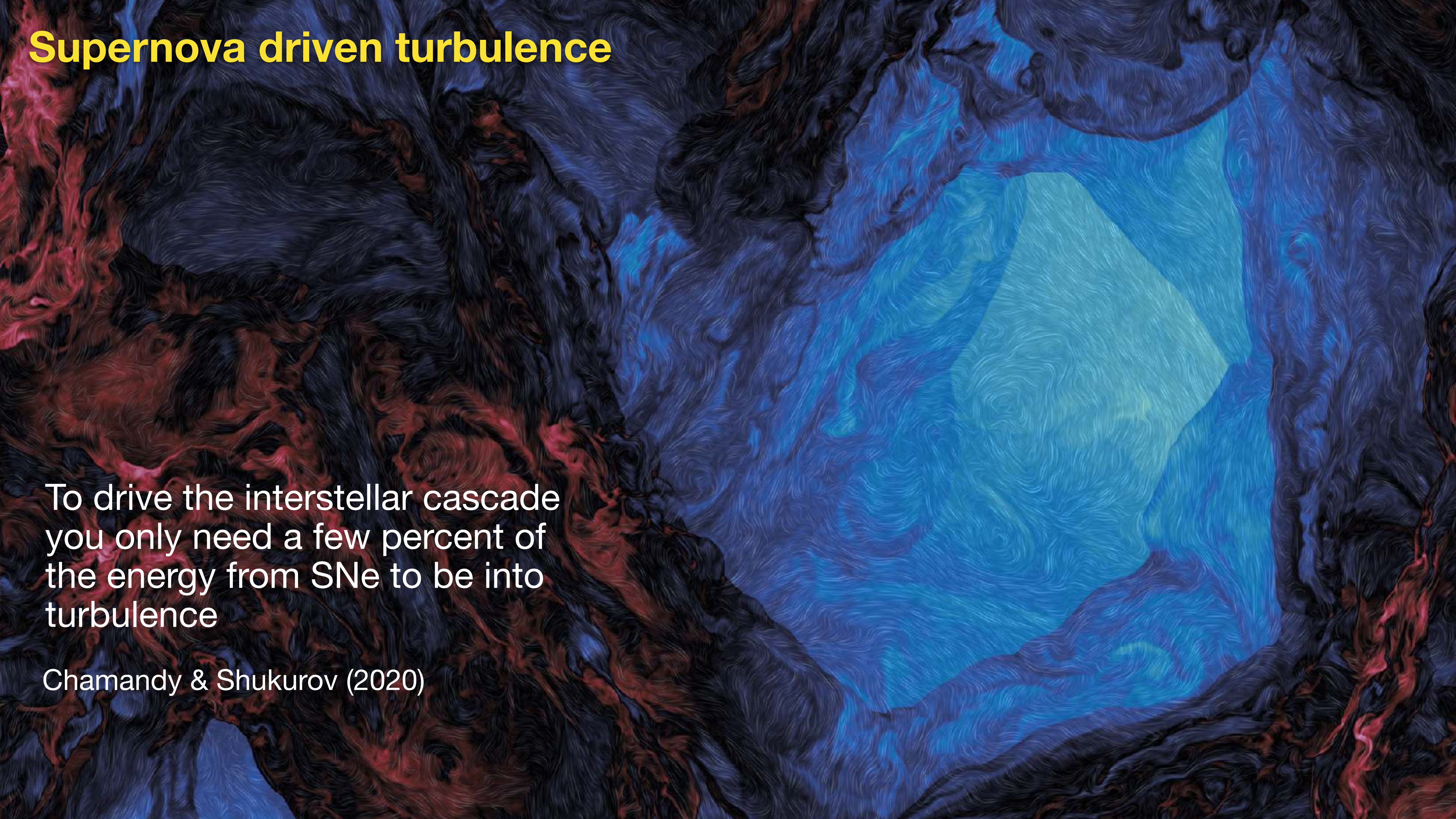
Plasma physicists

Supernova driven turbulence

To drive the interstellar cascade
you only need a few percent of
the energy from SNe to be into
turbulence

Chamandy & Shukurov (2020)

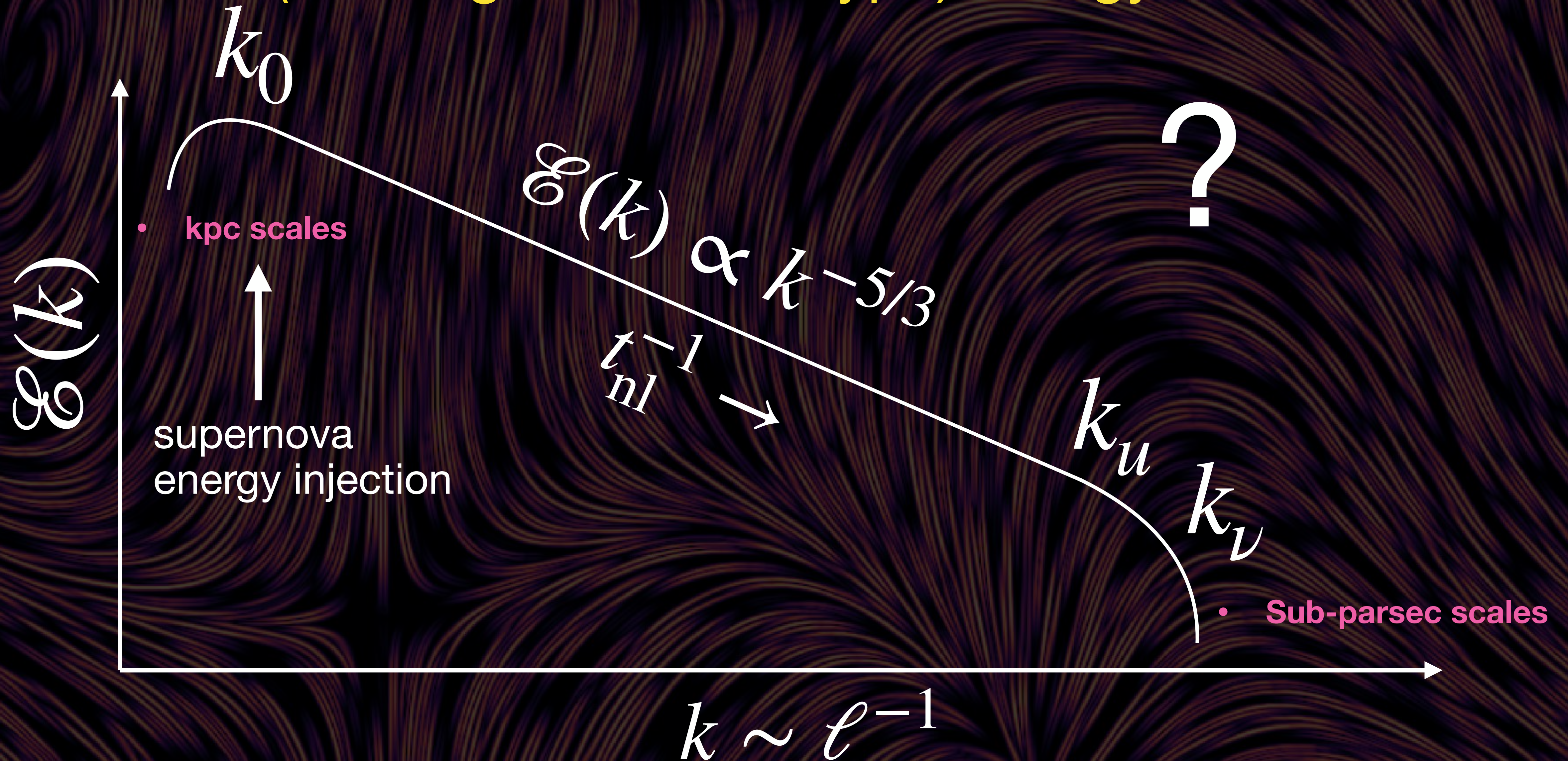
Supernova driven turbulence

The background image is a high-resolution simulation of a turbulent interstellar medium. It features a central, bright, irregularly shaped region in shades of light blue and white, which appears to be the remnant of a supernova explosion. This central region is surrounded by a complex, filamentary structure of dark blue and deep red, representing the turbulent gas and dust being driven by the explosion. The overall texture is highly detailed and swirling, characteristic of hydrodynamic simulations of turbulence.

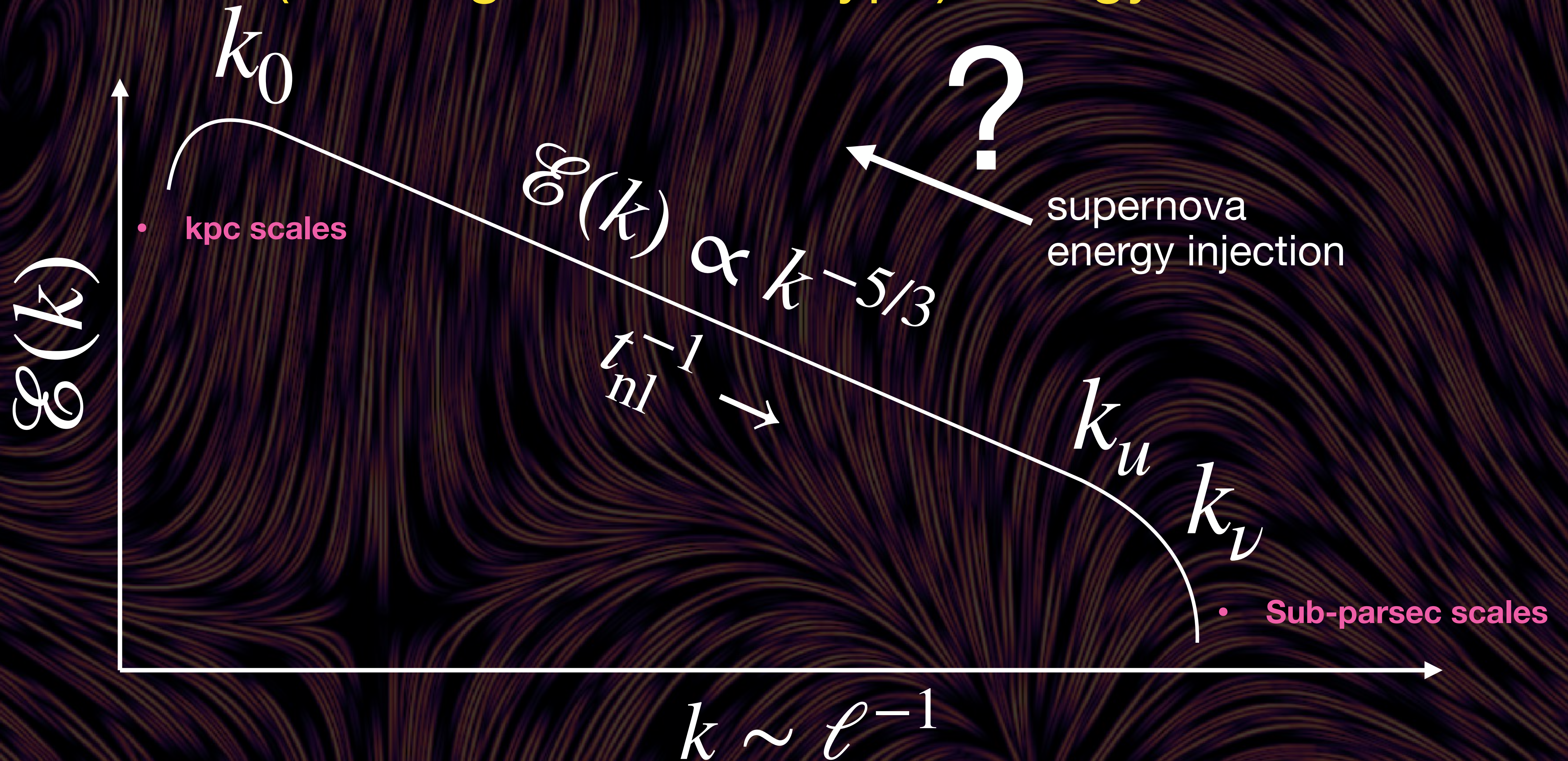
To drive the interstellar cascade
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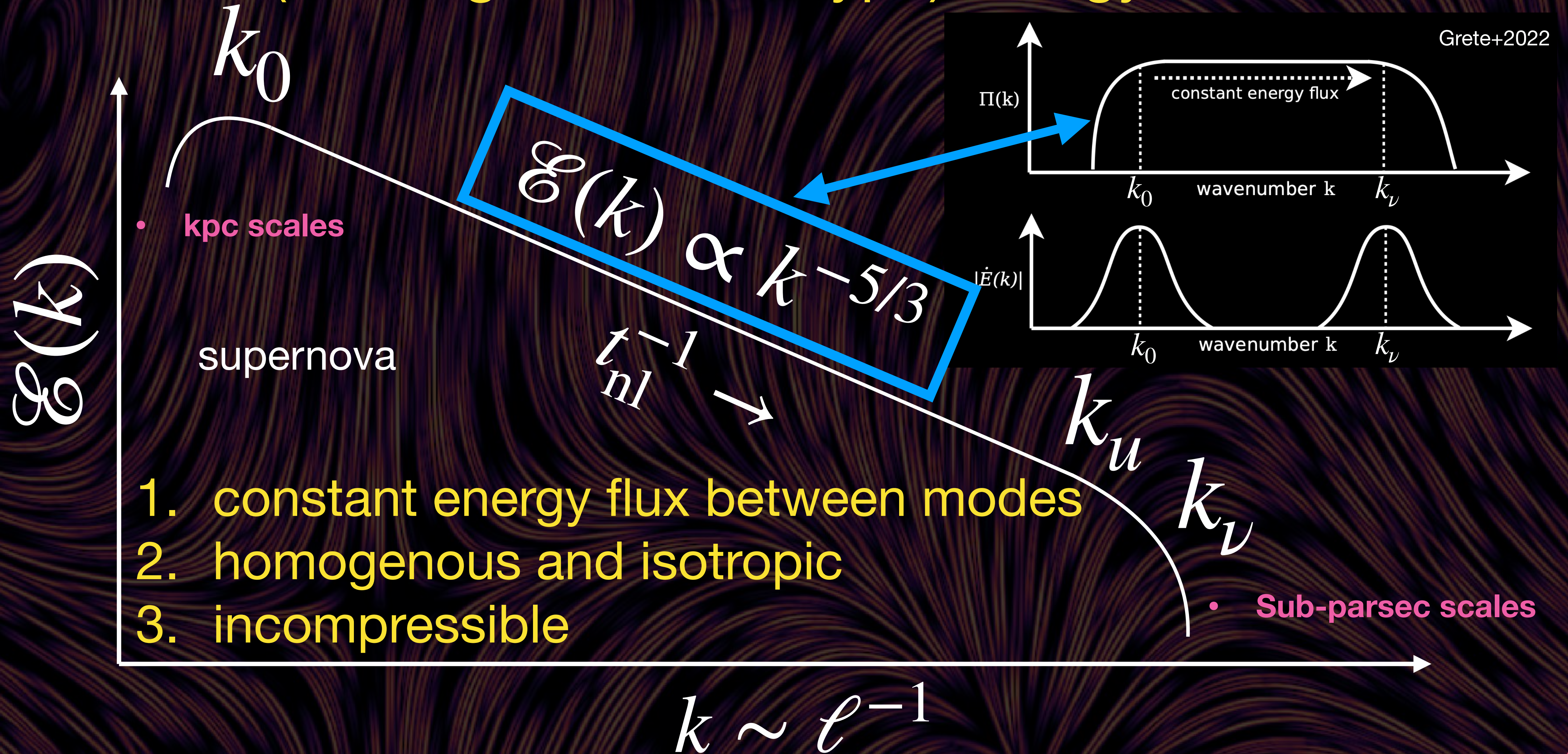
The (Kolmogorov, 1941-type) energy cascade



The (Kolmogorov, 1941-type) energy cascade

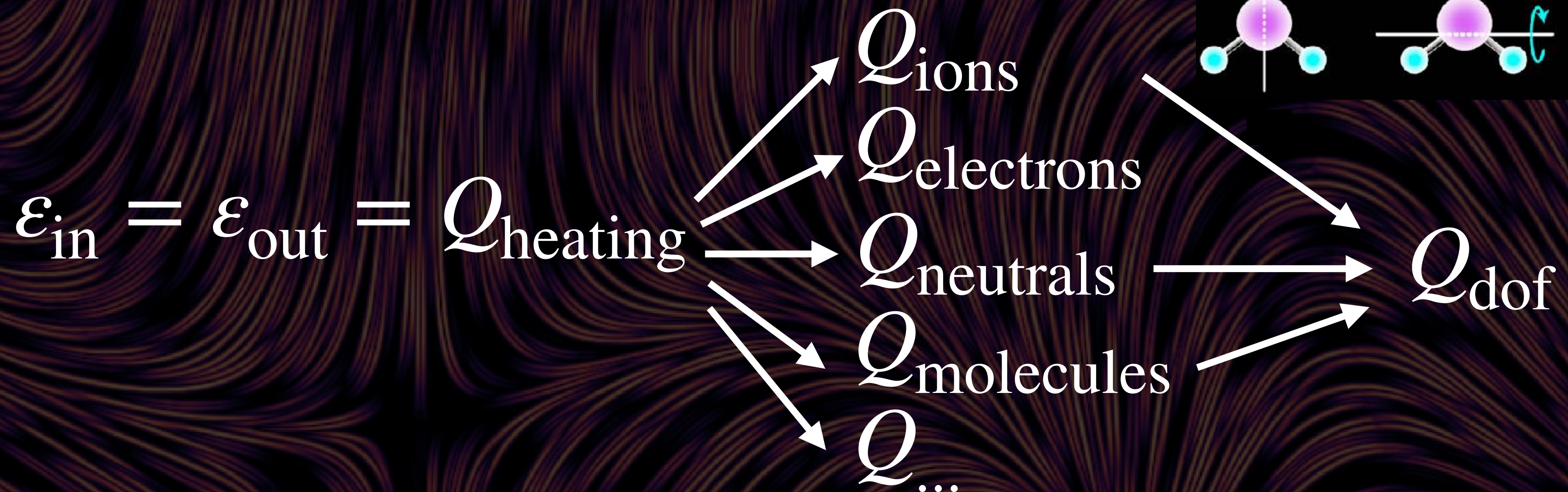


The (Kolmogorov, 1941-type) energy cascade



Energy flux density gives rise to Kolmogorov

$$\varepsilon \sim u^3/\ell \implies u \sim (\varepsilon\ell)^{1/3} \implies u^2(k) \sim k^{-5/3}$$



For details on heating and partition between ions and electrons and dof, read -> Greg Howes' latest work.
arxiv.org/abs/2402.12829

How to probe the energy flux density?

kinetic energy density

energy flux density from viscosity

$$\overbrace{\partial_t \mathcal{E}_{\text{kin}}} + \underbrace{\mathbf{u} \cdot \nabla \cdot \mathbb{F}_{\rho \mathbf{u}}}_{\text{energy flux density from transport}} = \frac{1}{\text{Re}} \overbrace{\mathbf{u} \cdot \mathbb{D}_\nu(\mathbf{u})}$$

How to probe the energy flux density?

kinetic energy density

energy flux density
from viscosity

$$\overbrace{\partial_t \mathcal{E}_{\text{kin}}} + \underbrace{\mathbf{u} \cdot \nabla \cdot \mathbb{F}_{\rho \mathbf{u}}}_{\text{energy flux density from transport}} = \frac{1}{\text{Re}} \overbrace{\mathbf{u} \cdot \mathbb{D}_\nu(\mathbf{u})}_{\text{energy flux density from viscosity}}$$

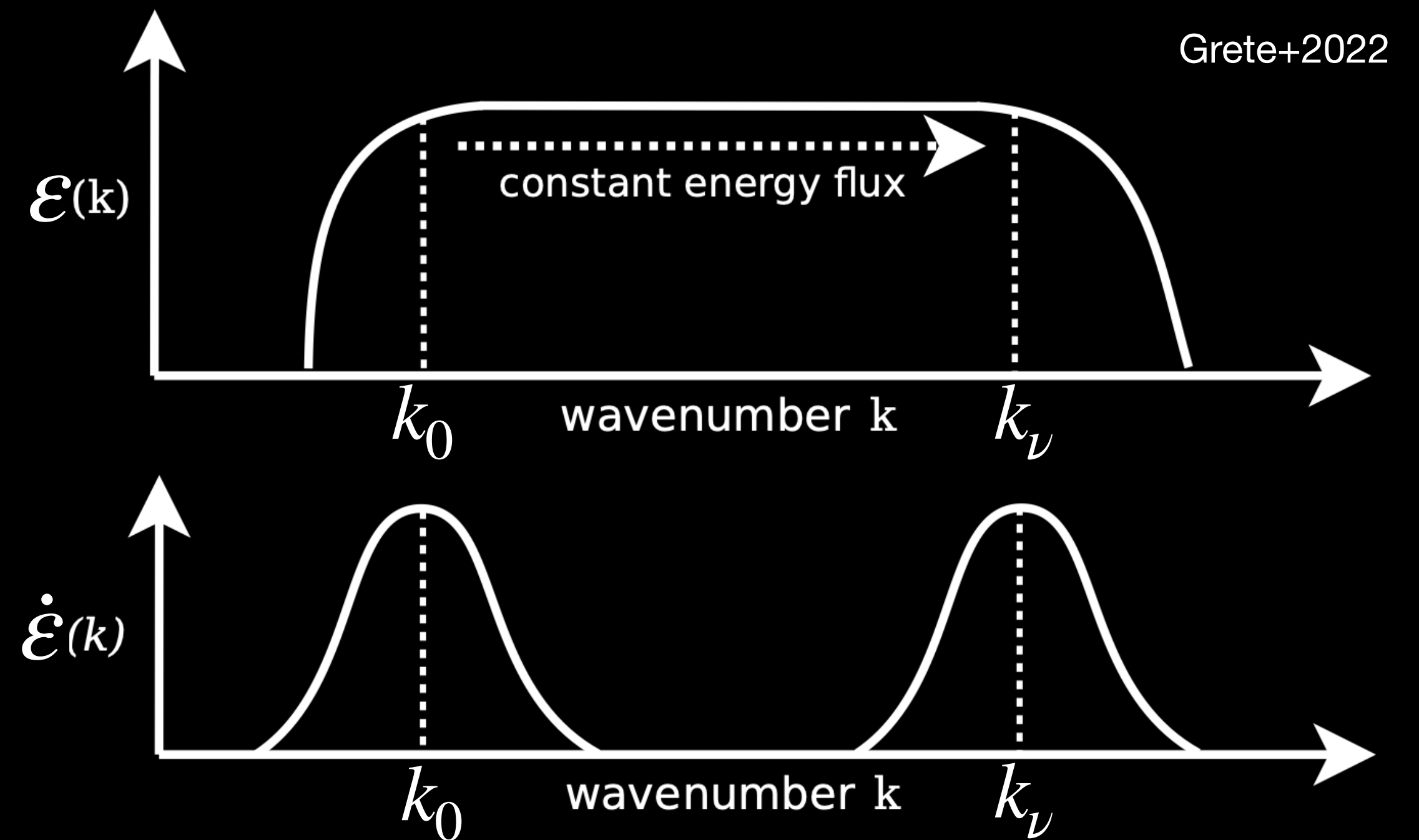
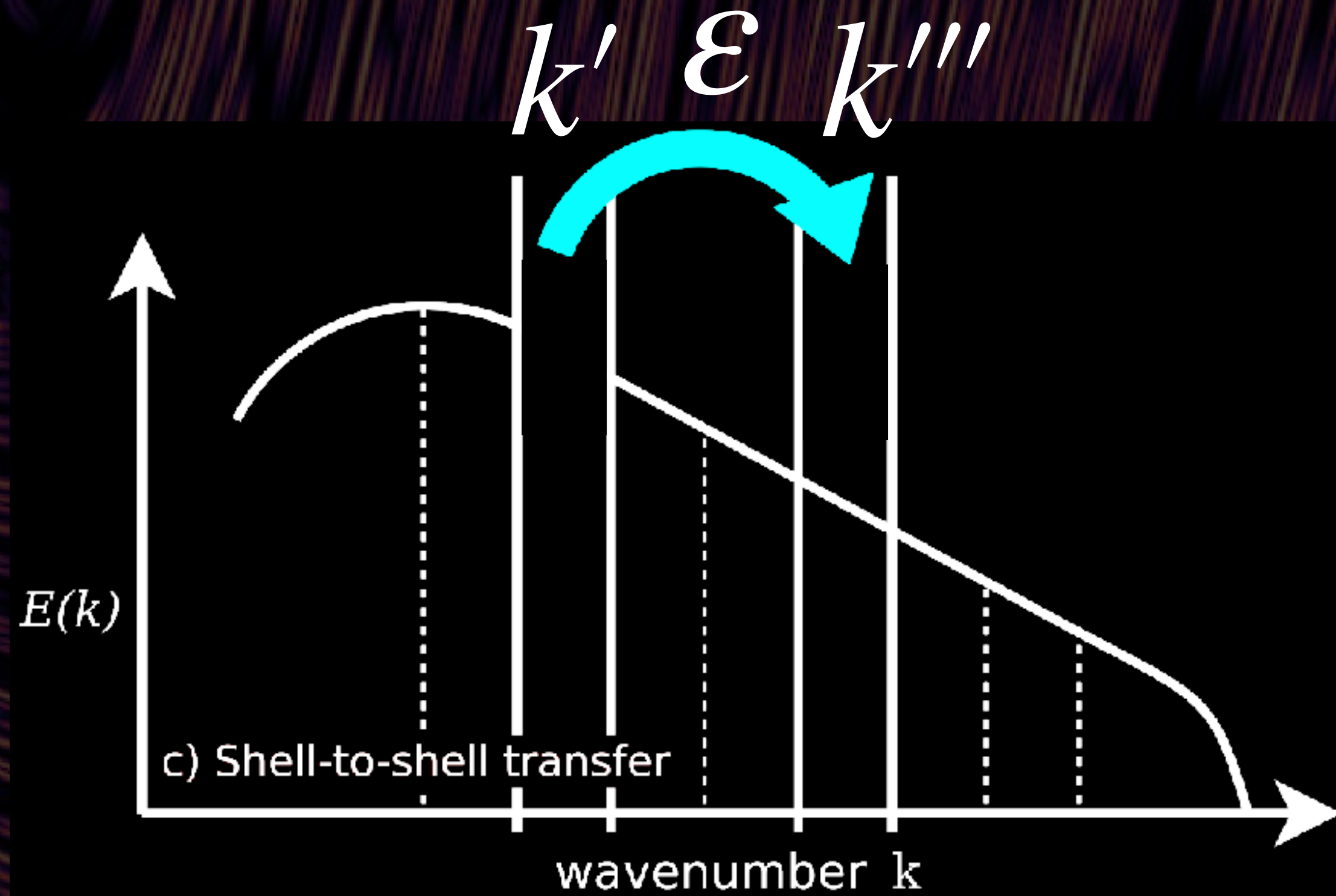
energy flux density
from transport

$$\mathbf{u} \cdot \nabla \cdot \mathbb{F}_{\rho \mathbf{u}} = \underbrace{\mathbf{u} \otimes \mathbf{u} : \nabla \otimes \mathbf{u}}_{\text{turbulence}} + \dots$$

That nonlinearity!

$$\mathbf{u} \cdot \mathbf{u} \cdot \nabla \otimes \mathbf{u}$$

But to understand the turbulence we must understand the flux from mode to mode!



Grete+2022a,b,23

Mininni+2005

But to understand the turbulence we must
understand the flux from mode to mode!

Momentum conservation:

$$\begin{array}{ccccc} \text{doner} & & \text{receiver} & & \\ \mathbf{k}' + \mathbf{k}'' + \mathbf{k}''' = 0 \\ \text{mediator} & & & & \end{array}$$

$$\begin{array}{ccccc} \text{doner} & & \text{receiver} & & \\ \mathbf{k}' \xrightarrow{\mathbf{k}''} \mathbf{k}''' = -\mathbf{k}''' \xrightarrow{\mathbf{k}''} \mathbf{k}' \\ \text{mediator} & & & & \end{array}$$

$$\mathbf{u}' = \mathbf{u}(\mathbf{r}') = \int \delta^3(\mathbf{k} - \mathbf{k}') \mathbf{u}(\mathbf{k}) \exp \{2\pi i \mathbf{k} \cdot \mathbf{r}\}$$

But to understand the turbulence we must
understand the flux from mode to mode!

kinetic energy density

$$\frac{1}{2} \overbrace{\partial_t \rho \mathbf{u}' \cdot \mathbf{u}'''} + \mathbf{u}''' \cdot \underbrace{\nabla \cdot \mathbb{F}_{\rho \mathbf{u}}(\mathbf{u}'', \mathbf{u}')}_{\text{energy flux density}} = 0$$

$$\mathbf{k}' \xrightarrow{\mathbf{k}''} \mathbf{k}'''$$

energy flux density
from transport between $\mathbf{u}'$ and $\mathbf{u}'''$

$$\mathbf{u}''' \cdot \nabla \cdot \mathbb{F}(\mathbf{u}'', \mathbf{u}')_{\rho \mathbf{u}} = \mathbf{u}''' \otimes \mathbf{u}'' : \nabla \otimes \mathbf{u}' + \dots$$

$$\varepsilon \sim u^3 / \ell$$

But to understand the turbulence we must
understand the flux from mode to mode!

$$\mathbf{u}''' \cdot \nabla \cdot \mathbb{F}(\mathbf{u}'', \mathbf{u}')_{\rho \mathbf{u}} = \mathbf{u}''' \otimes \mathbf{u}'' : \nabla \otimes \mathbf{u}' + \dots$$

$$\mathcal{T}_{uu}(k', k''' | k'') = - \int dV \mathbf{u}''' \otimes \mathbf{u}'' : \nabla \otimes \mathbf{u}'$$

$\mathbf{k}' \xrightarrow{\mathbf{k}''} \mathbf{k}'''$

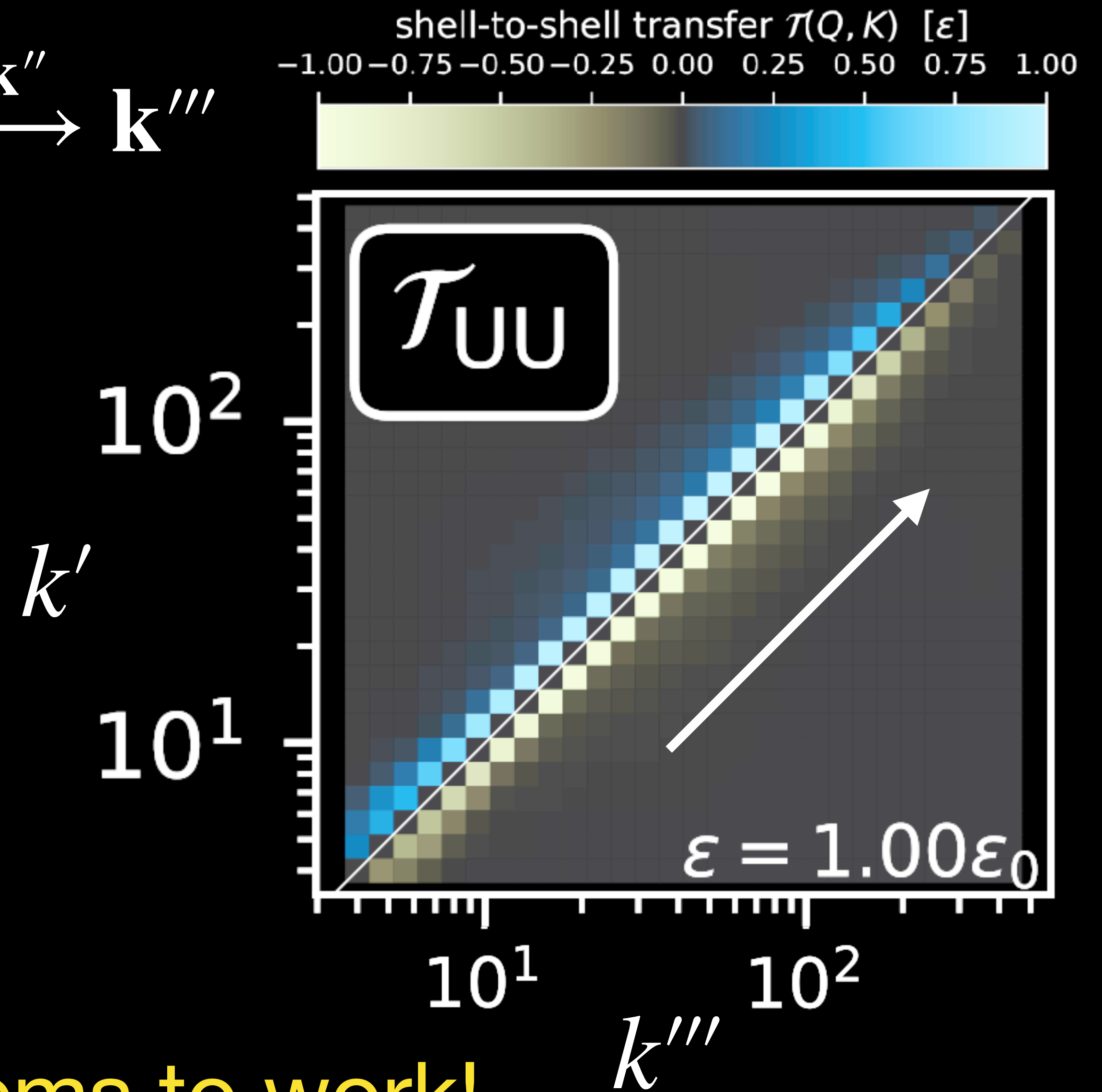
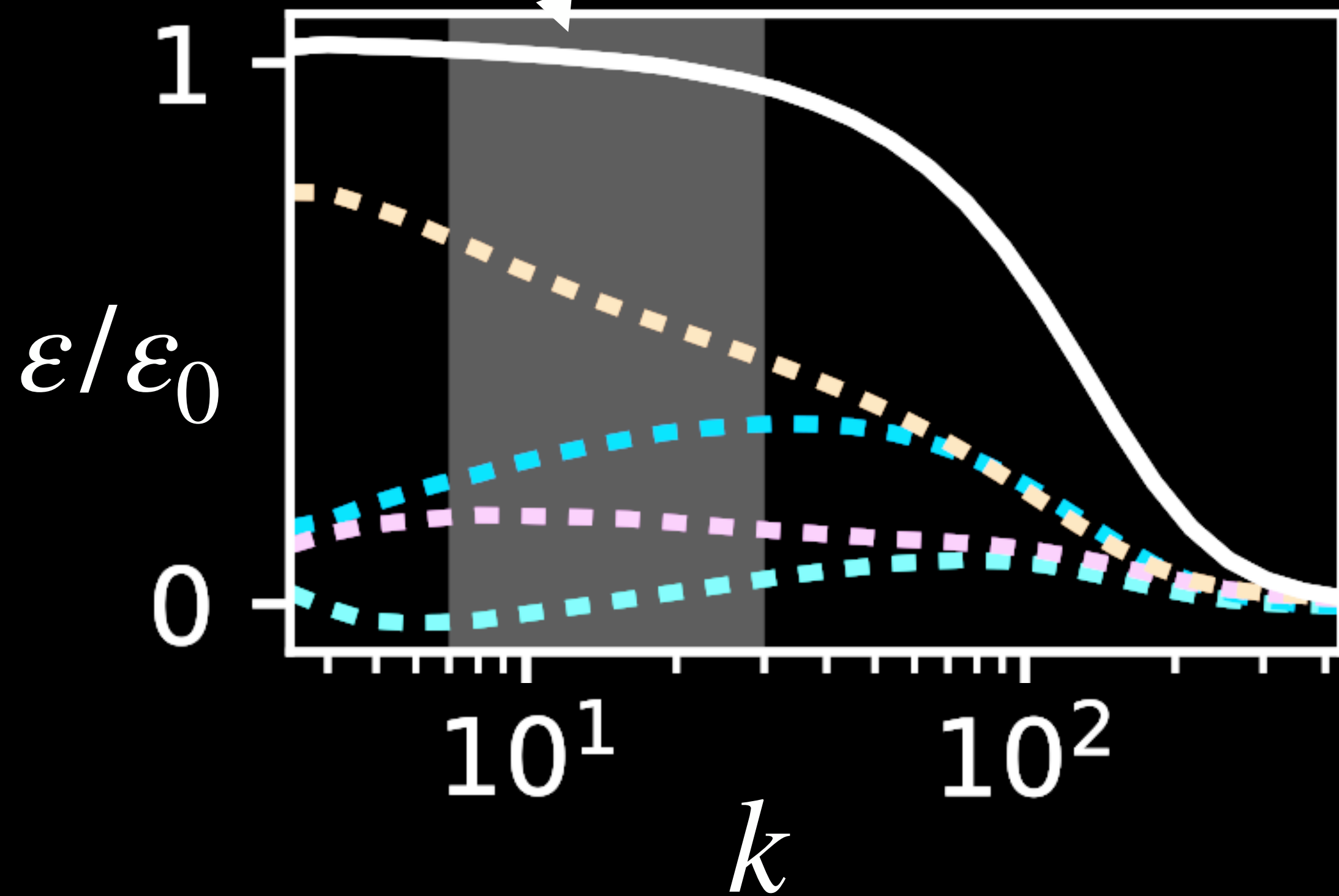
$$\varepsilon \sim u^3 / \ell$$

Compressible MHD turbulence in a box

Grete+2022

$$\mathbf{k}' \xrightarrow{\mathbf{k}''} \mathbf{k}'''$$

$$\varepsilon \sim u^3/\ell \approx \text{const.}$$



Kolmogorov assumption seems to work!

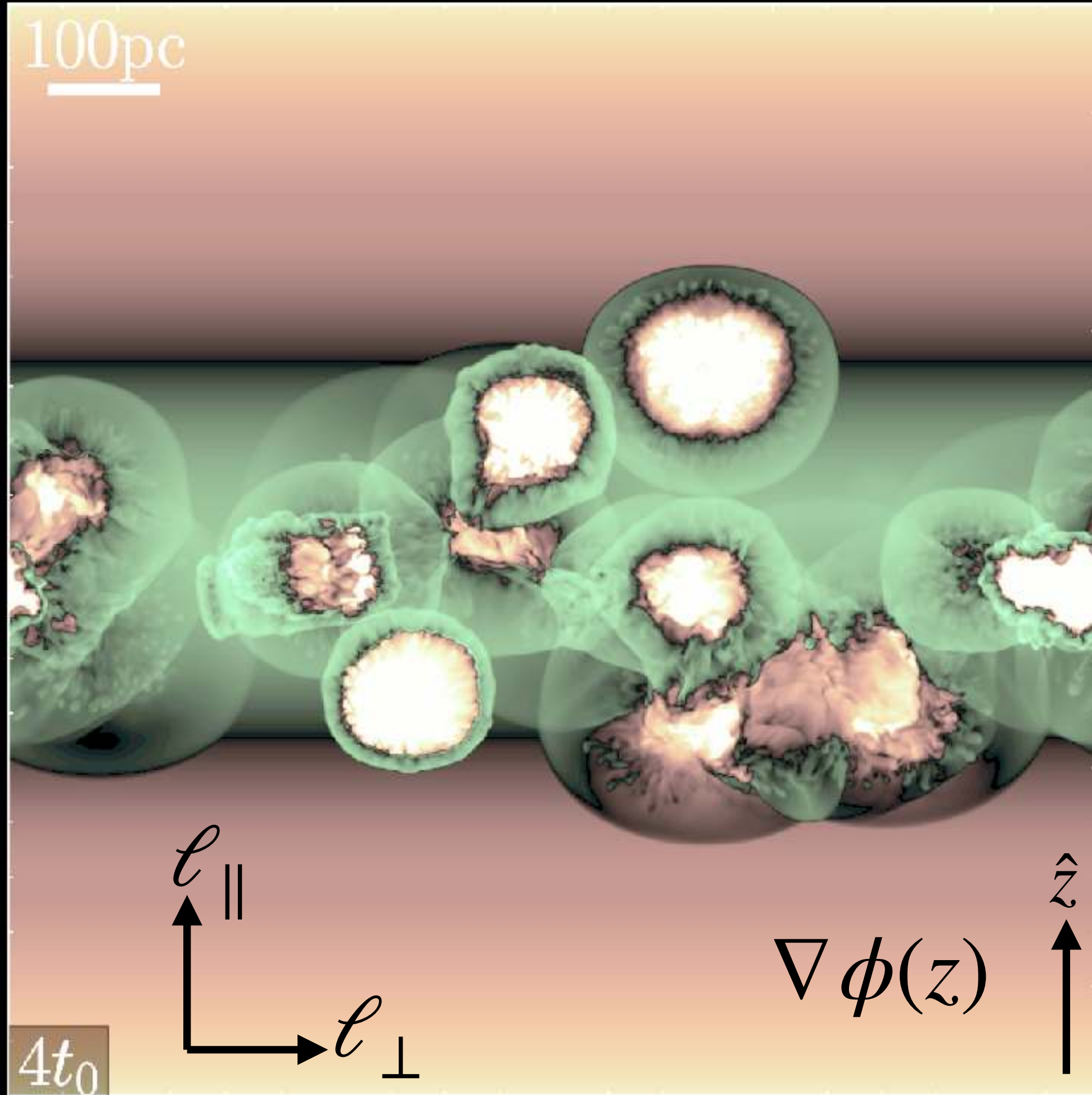
Let's do this for more realistic ISM turbulence

$L = 1 \text{ kpc}$

Martizzi+2016

RAMSES (Teyssier 2002)

Supernova driven
gravito-hydro dynamical model



10^2

10^1

10^0

10^{-1}

10^{-2}

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = \dot{n}_{\text{SNe}} M_{\text{ej}}, \quad (1)$$

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u} + P \mathbb{I}) = -\rho \nabla \phi + \dot{n}_{\text{SNe}} \mathbf{p}_{\text{SNe}}(Z, n_{\text{H}}), \quad (2)$$

$$\frac{\partial \rho e}{\partial t} + \nabla \cdot [\rho (e + P) \mathbf{u}] = -n_{\text{H}}^2 \Lambda - \quad (3)$$

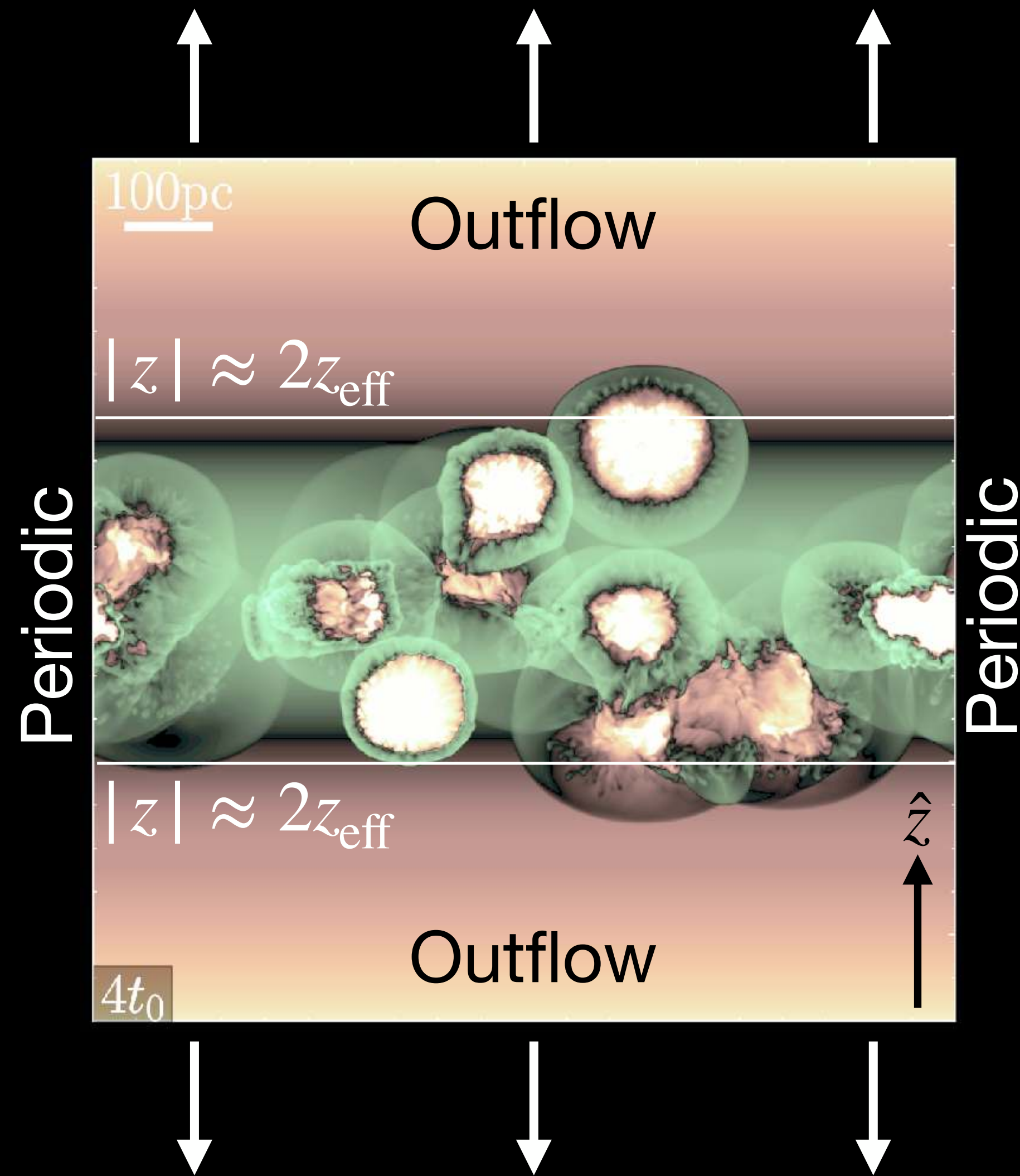
$$\rho \mathbf{u} \cdot \nabla \phi + \dot{n}_{\text{SNe}} \left[E_{\text{th,SNe}}(Z, n_{\text{H}}) + \frac{\mathbf{p}_{\text{SNe}}(Z, n_{\text{H}})^2}{2(M_{\text{ej}} + M_{\text{swept}})} \right], \quad (4)$$

$$e = \epsilon + \frac{u^2}{2}, \quad P = \frac{2}{3} \rho e, \quad (5)$$

$$N_{\text{grid}}^3 = 1024^3 \implies \Delta x \sim 1 \text{ pc}$$

$$\text{numerical diss.} \implies \Delta x \sim 10 \text{ pc}$$

Let's do this for more realistic ISM turbulence



Static gravitational potential

$$\phi(z) = \underbrace{2\pi G \Sigma_* \left(\sqrt{z^2 - z_0^2} - z_0 \right)}_{\text{stratified disk}} + \underbrace{\frac{2\pi G \rho_{\text{halo}}}{3} z^2}_{\text{spherical halo}}$$

Supernova driving prescription

$$\dot{n}_{\text{SNe}} = \frac{\dot{\Sigma}_*}{2z_{\text{eff}} 100 M_{\odot}} \quad \text{1 SNe per } 100 M_{\odot} \text{ of SF}$$

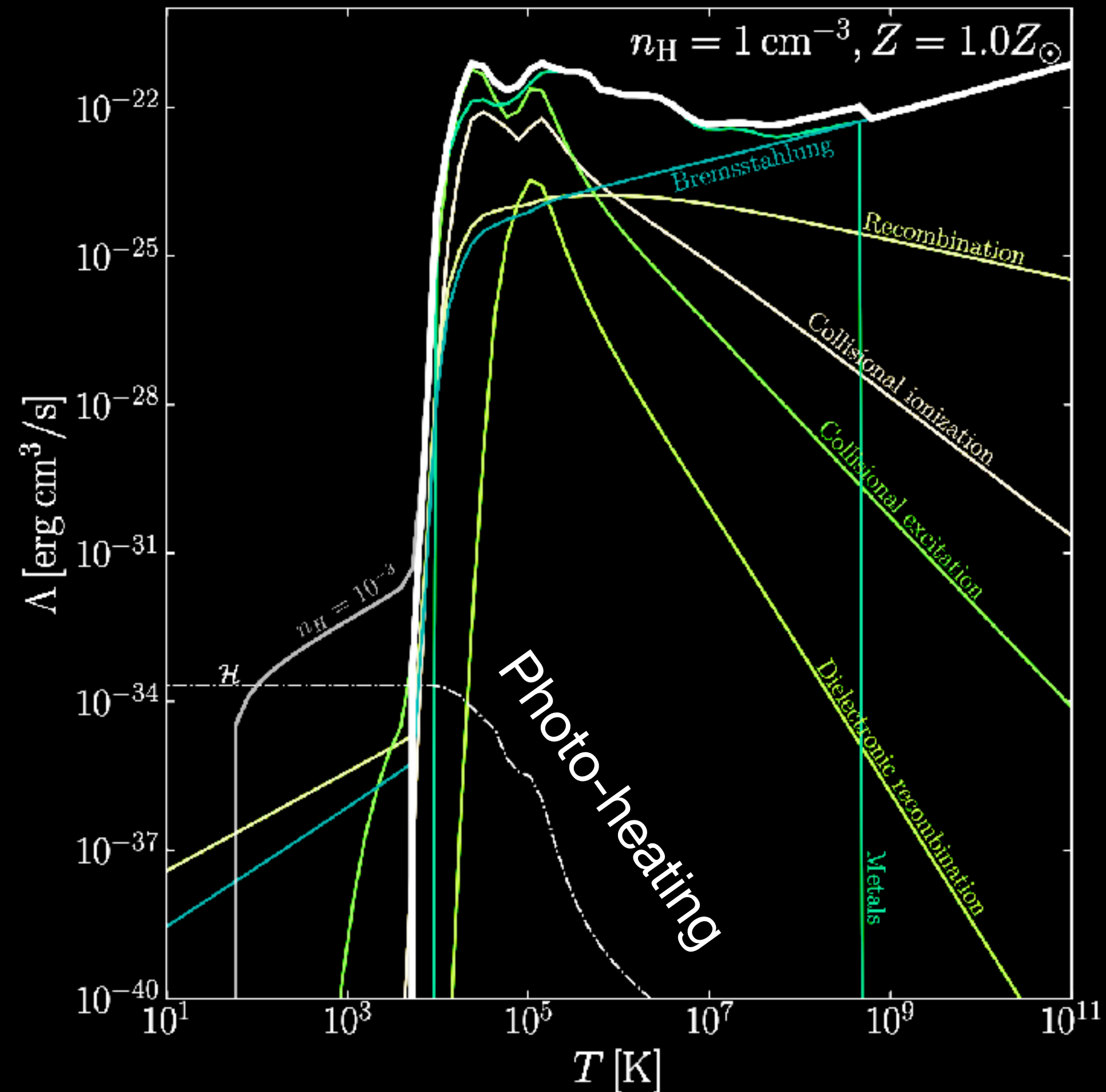
$$\dot{\Sigma}_* \propto \Sigma^{1.4} \quad \text{KS relation}$$

$$p(z) = \begin{cases} \frac{1}{4z_{\text{eff}}}, & |z| \leq 2z_{\text{eff}} \\ 0, & |z| > 2z_{\text{eff}} \end{cases}$$

Inject many passive scalar metals into medium

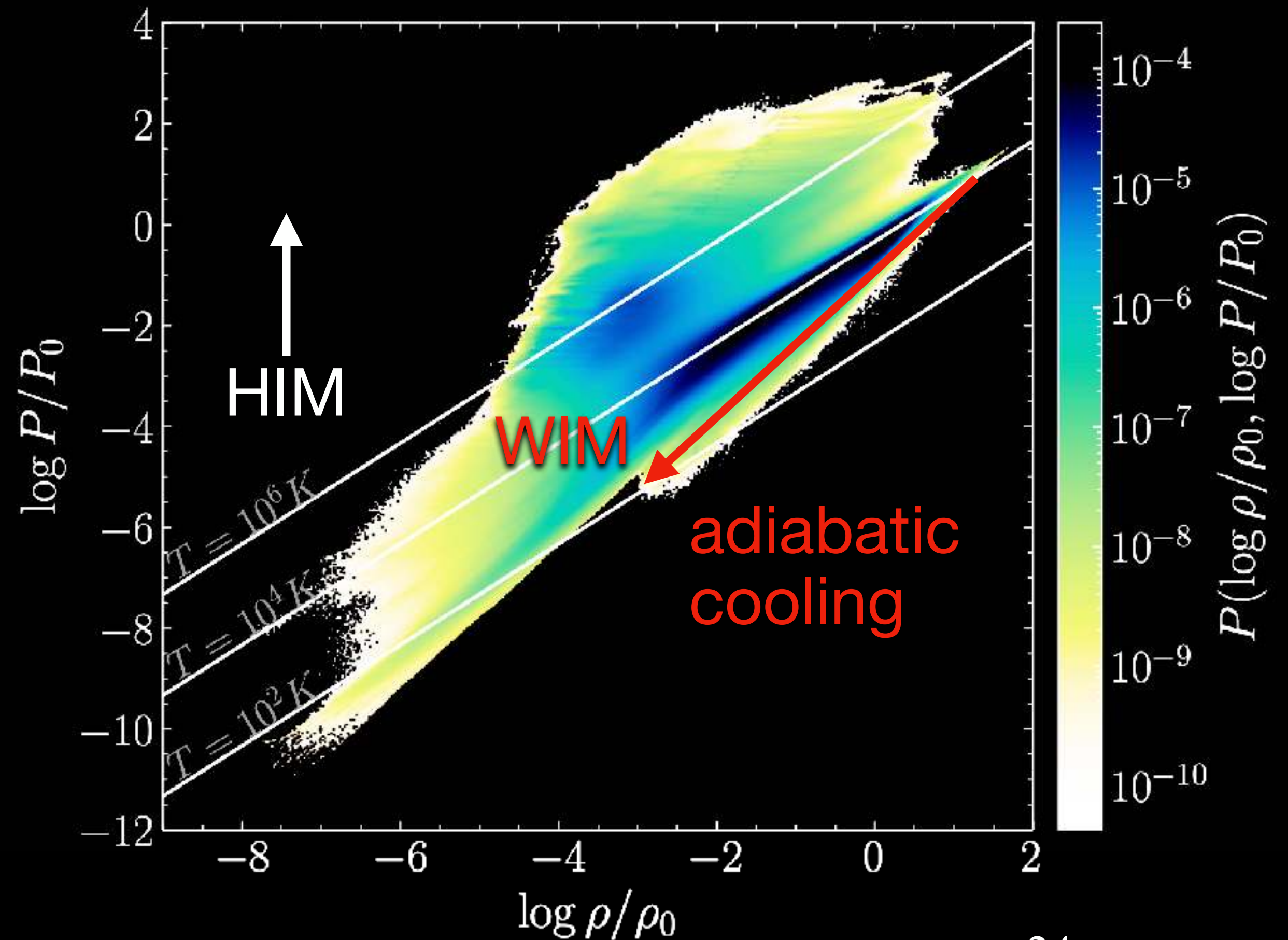
Let's do this for more realistic ISM turbulence

Cooling function Theuns+(1998)
Sutherland & Dopita(1993)



HI, HII, HeI, HeII, HeIII and free electrons

Phase space



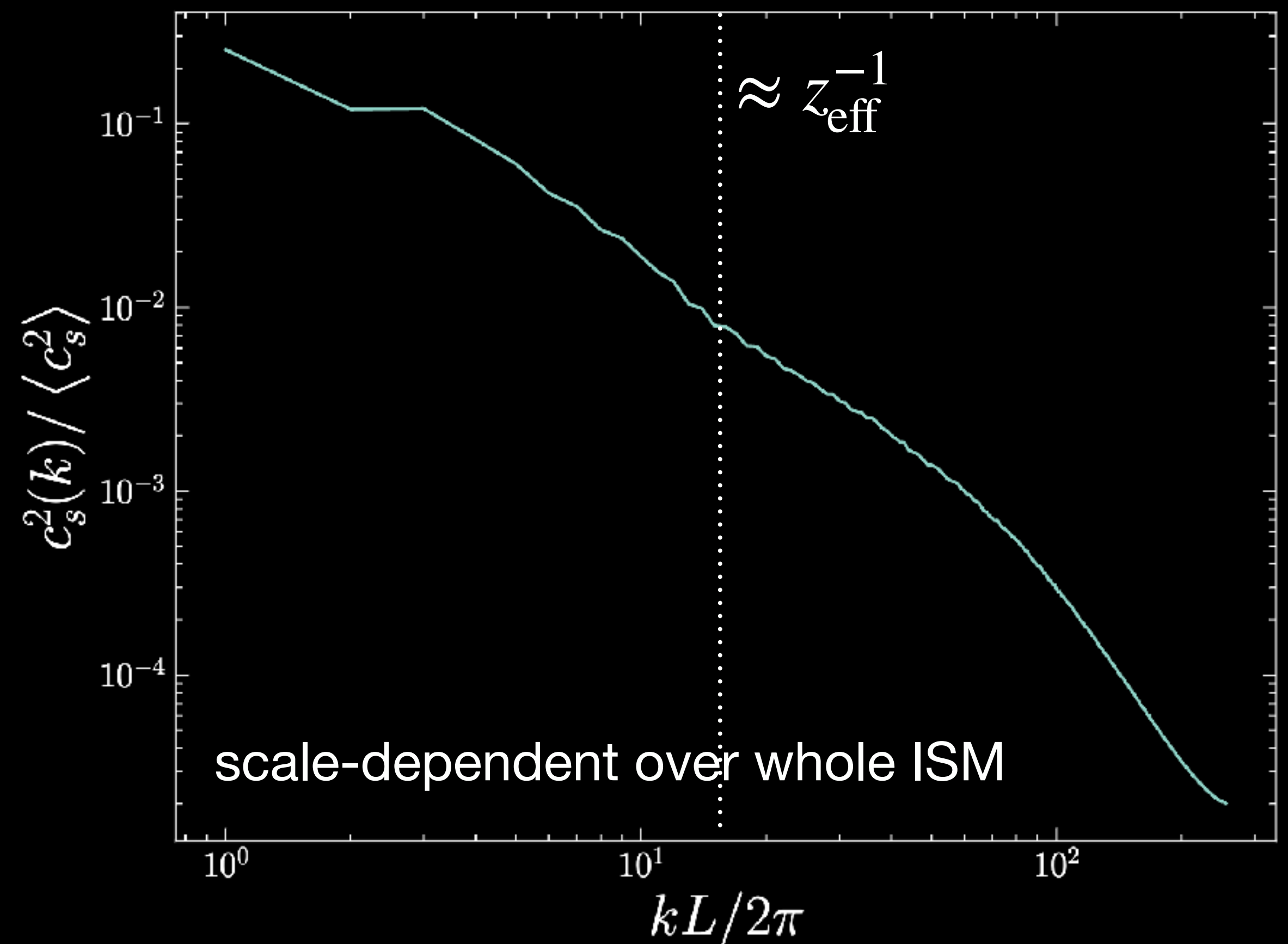
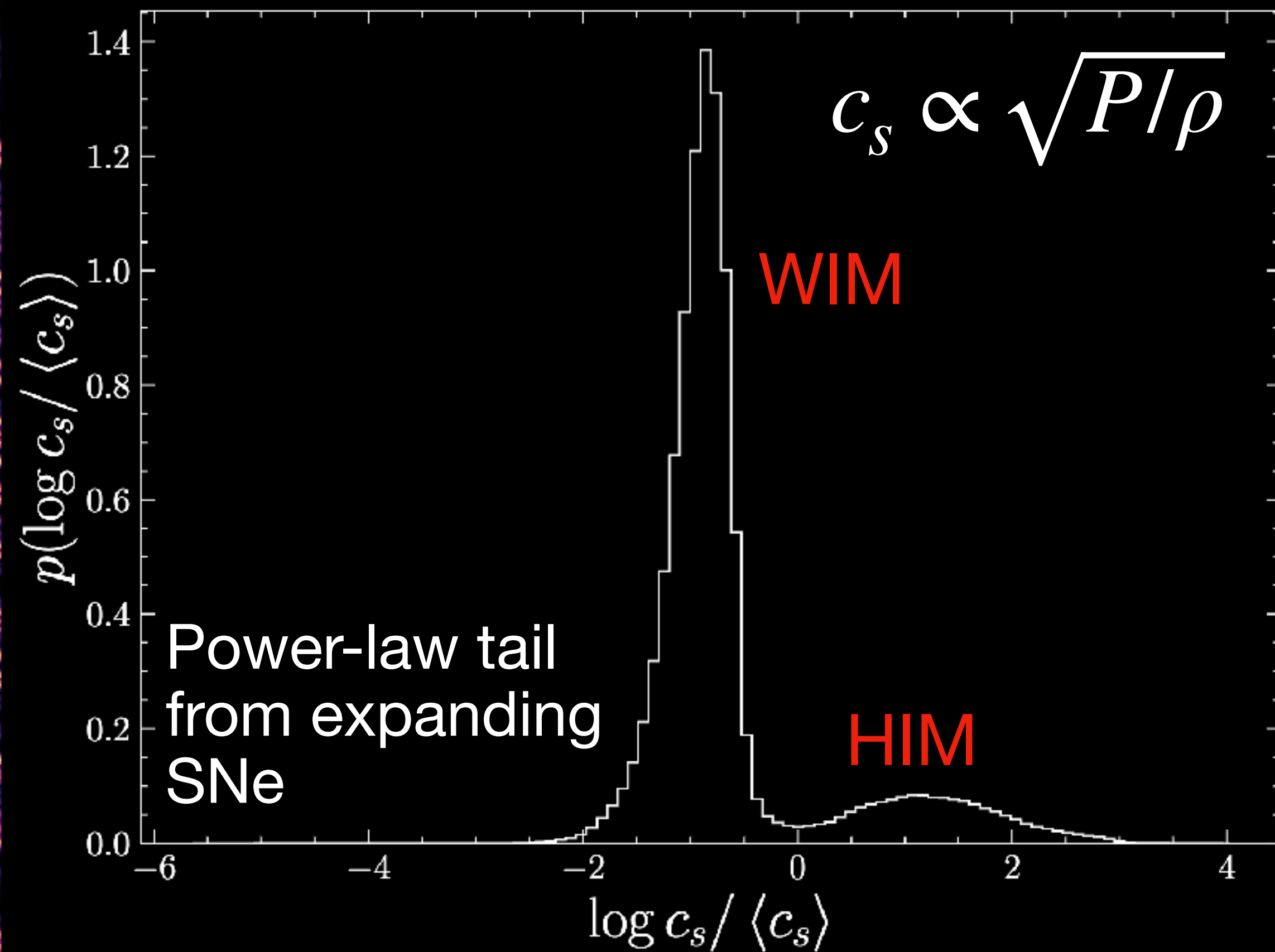
$$\rho_0 = 2.1 \times 10^{-24} \text{ g cm}^{-3}$$

$$P_0 = 2.2 \times 10^{-12} \text{ Ba}$$

Let's do this for more realistic ISM turbulence

Sound speed statistics

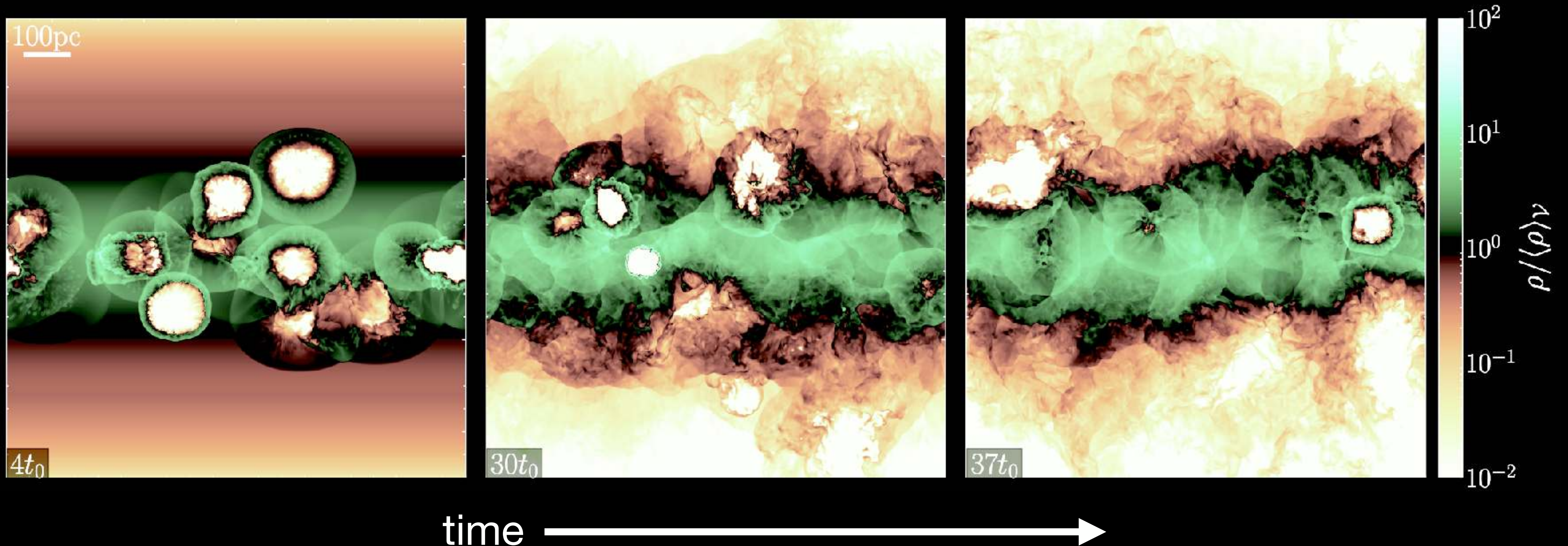
$$N_{\text{grid}}^3 = 512^3$$



Embracing the idea of a turbulent ISM is to embrace that every quantity is scale-dependent!

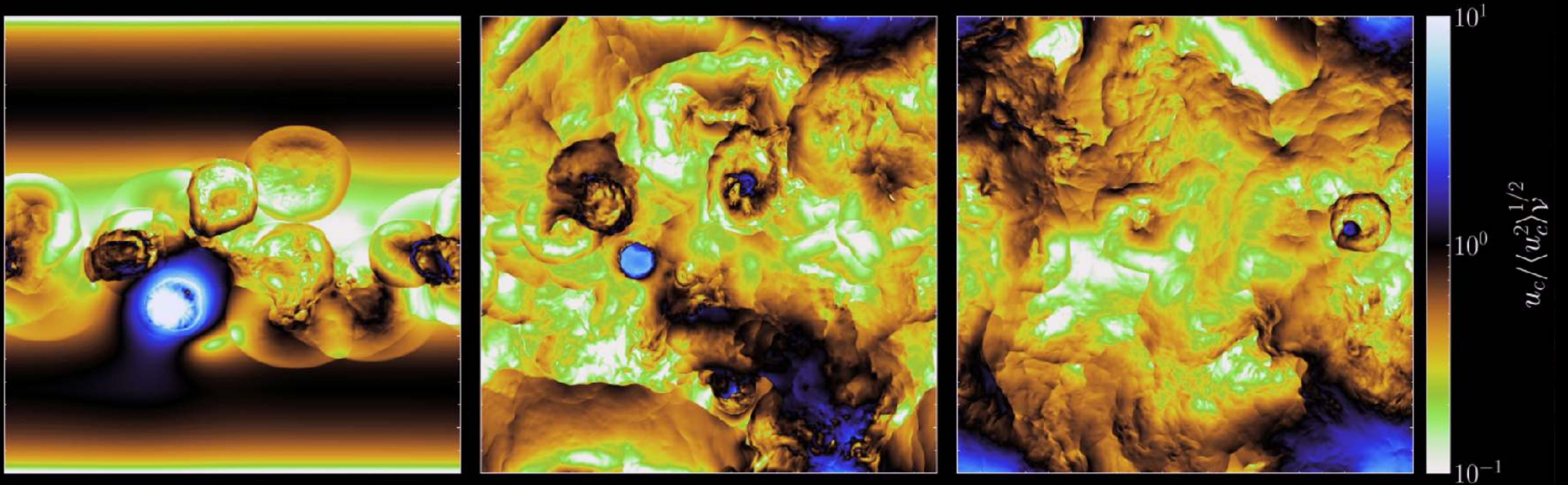
Let's do this for more realistic ISM turbulence

The journey towards stationarity: the mass density



Let's do this for more realistic ISM turbulence

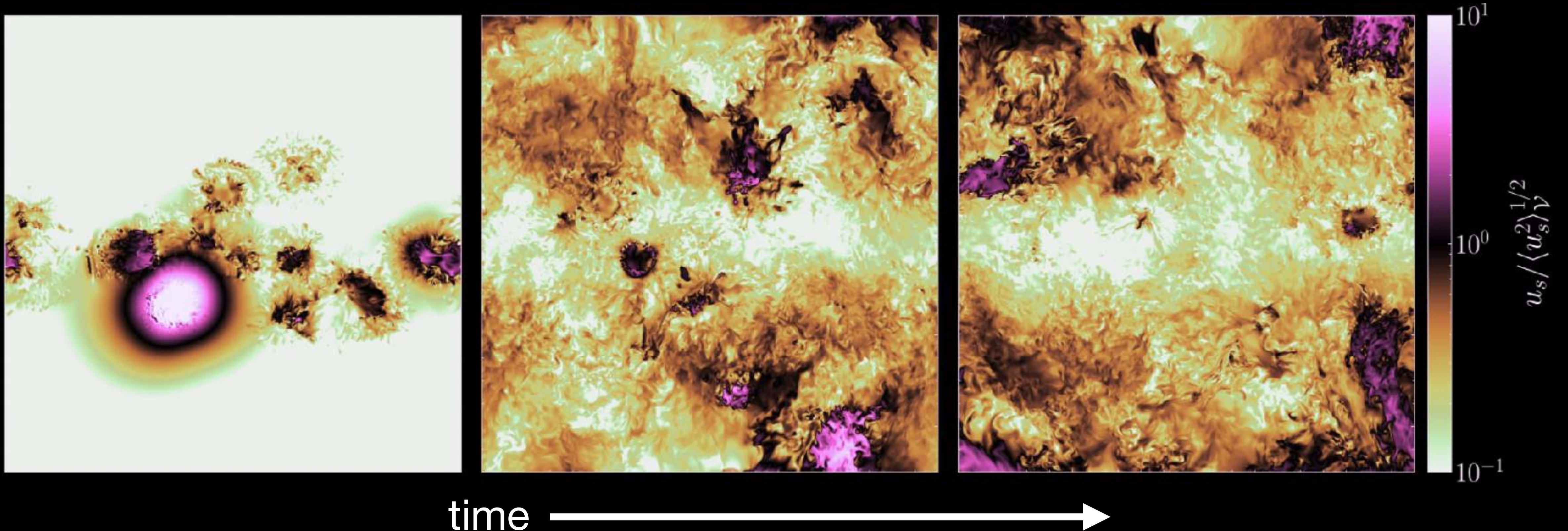
The journey towards stationarity: compressible modes



the velocity modes don't seem to care about the phases...

Let's do this for more realistic ISM turbulence

The journey towards stationarity: solenoidal modes

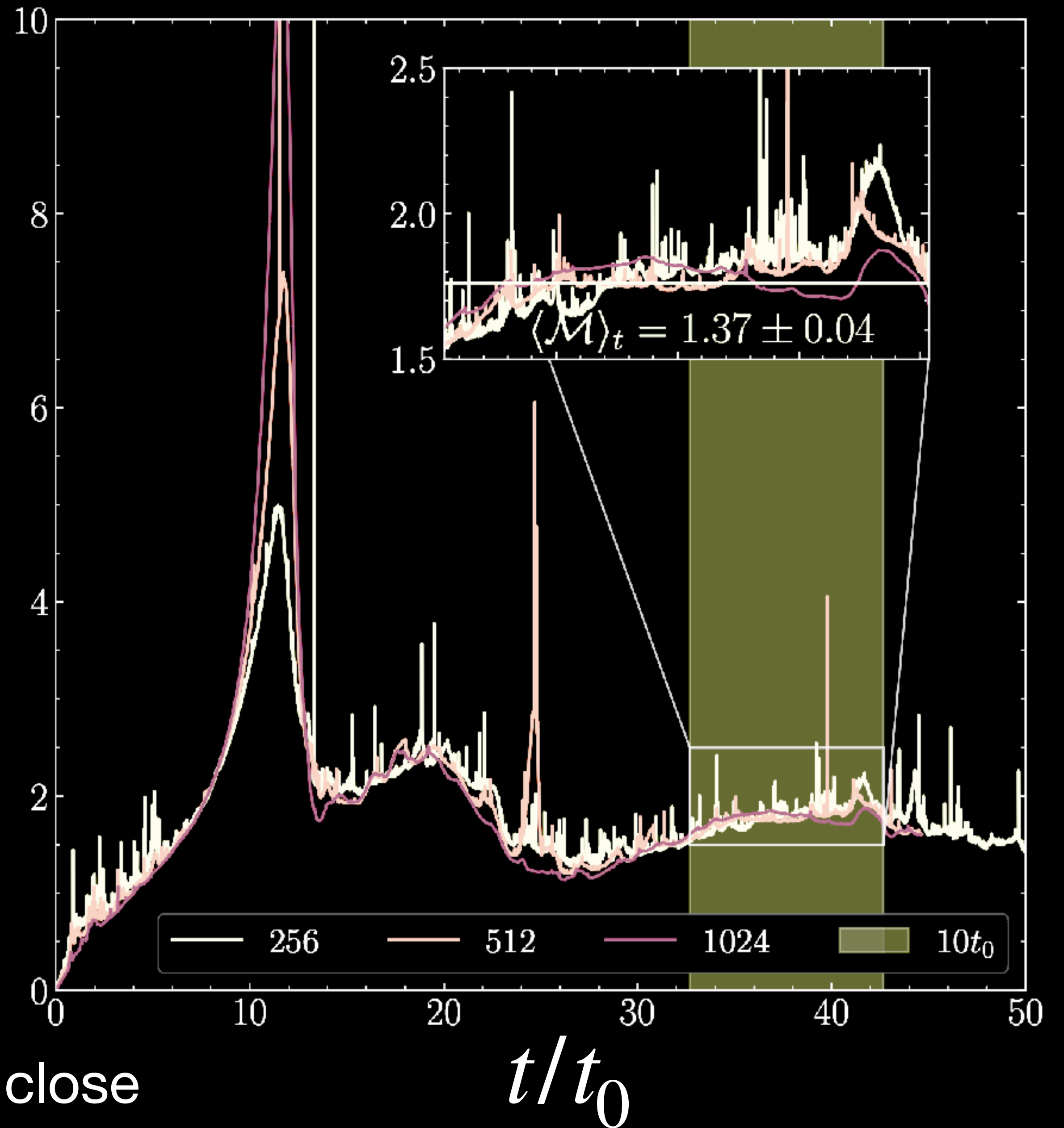


Stationarity and Mach number

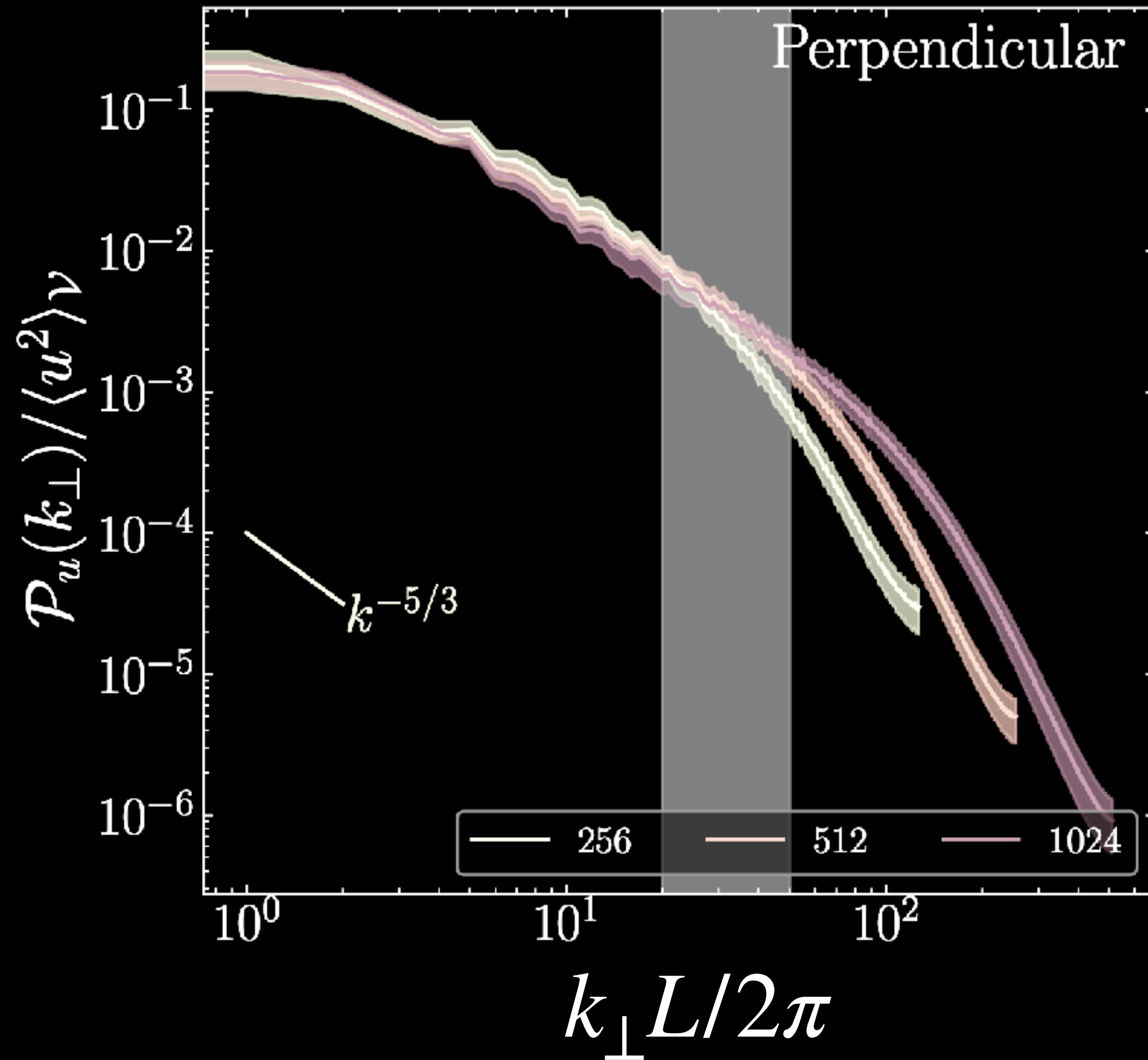
$$\mathcal{M} = \left\langle \left(\frac{u}{c_s} \right)^2 \right\rangle^{1/2}$$

$$t_0 \sim \frac{z_{\text{eff}}}{\langle u^2 \rangle^{1/2}} \sim 80 \text{ Myr}$$

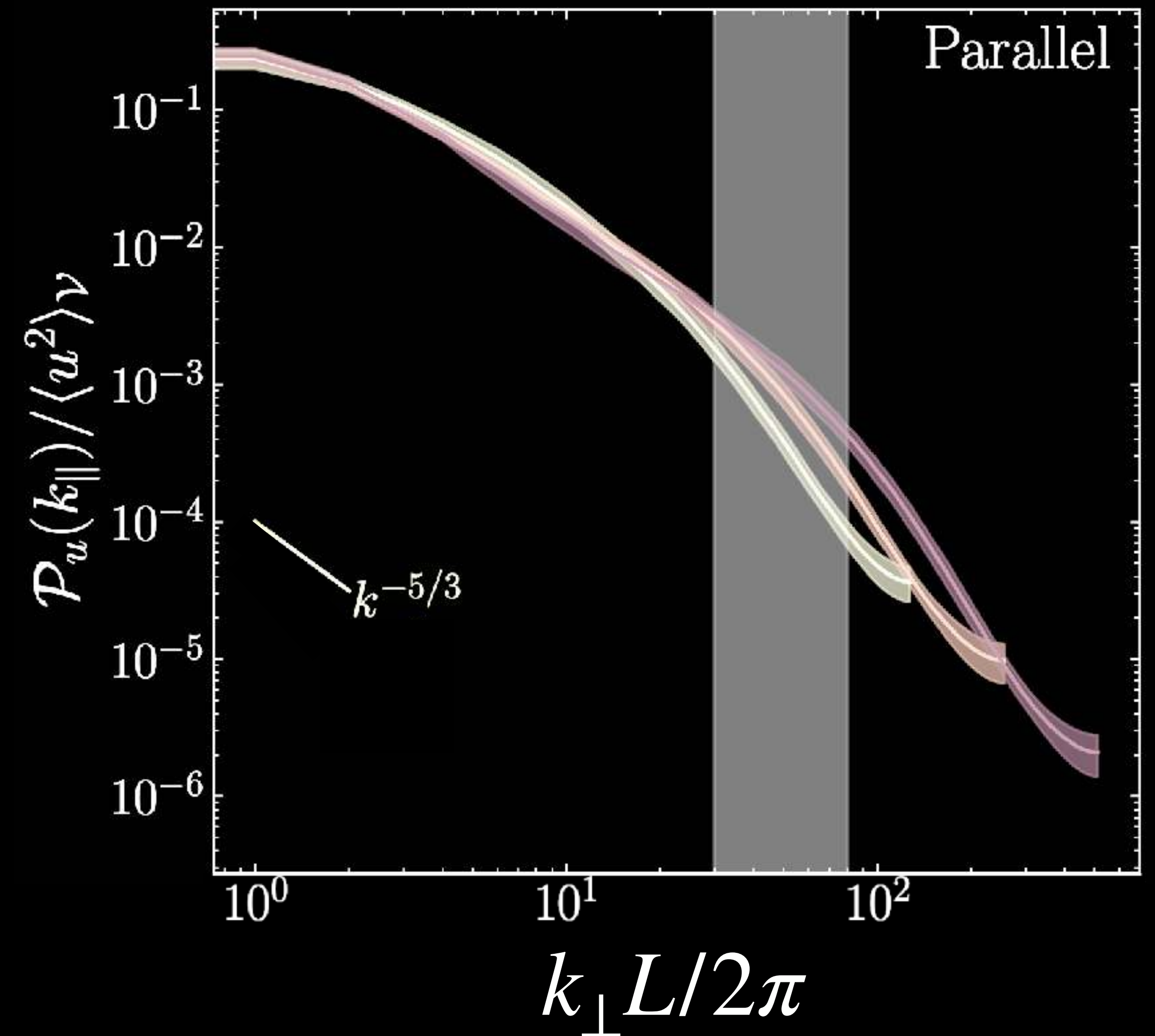
Reaches a close to stationary state... close



Velocity Spectra

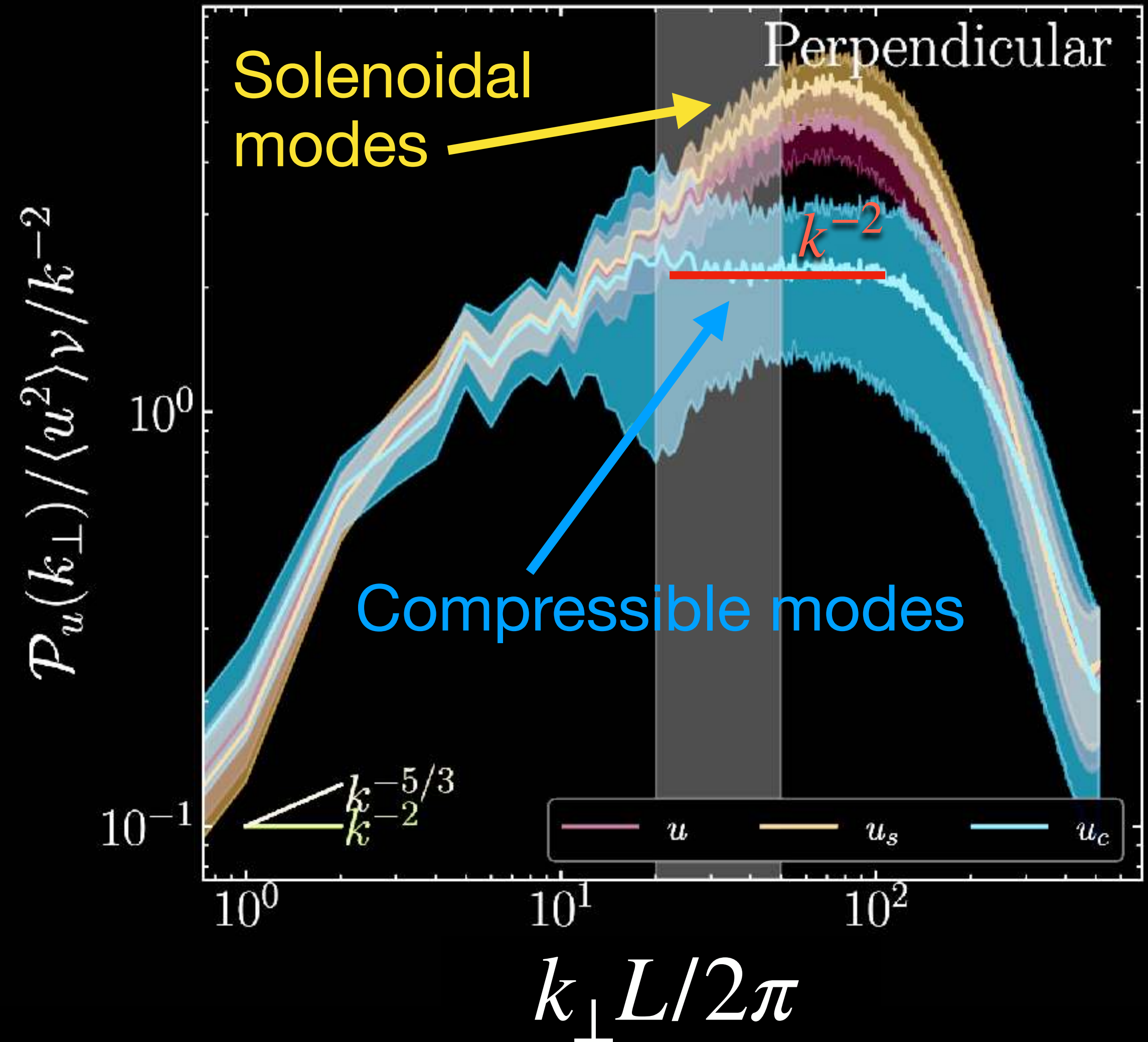
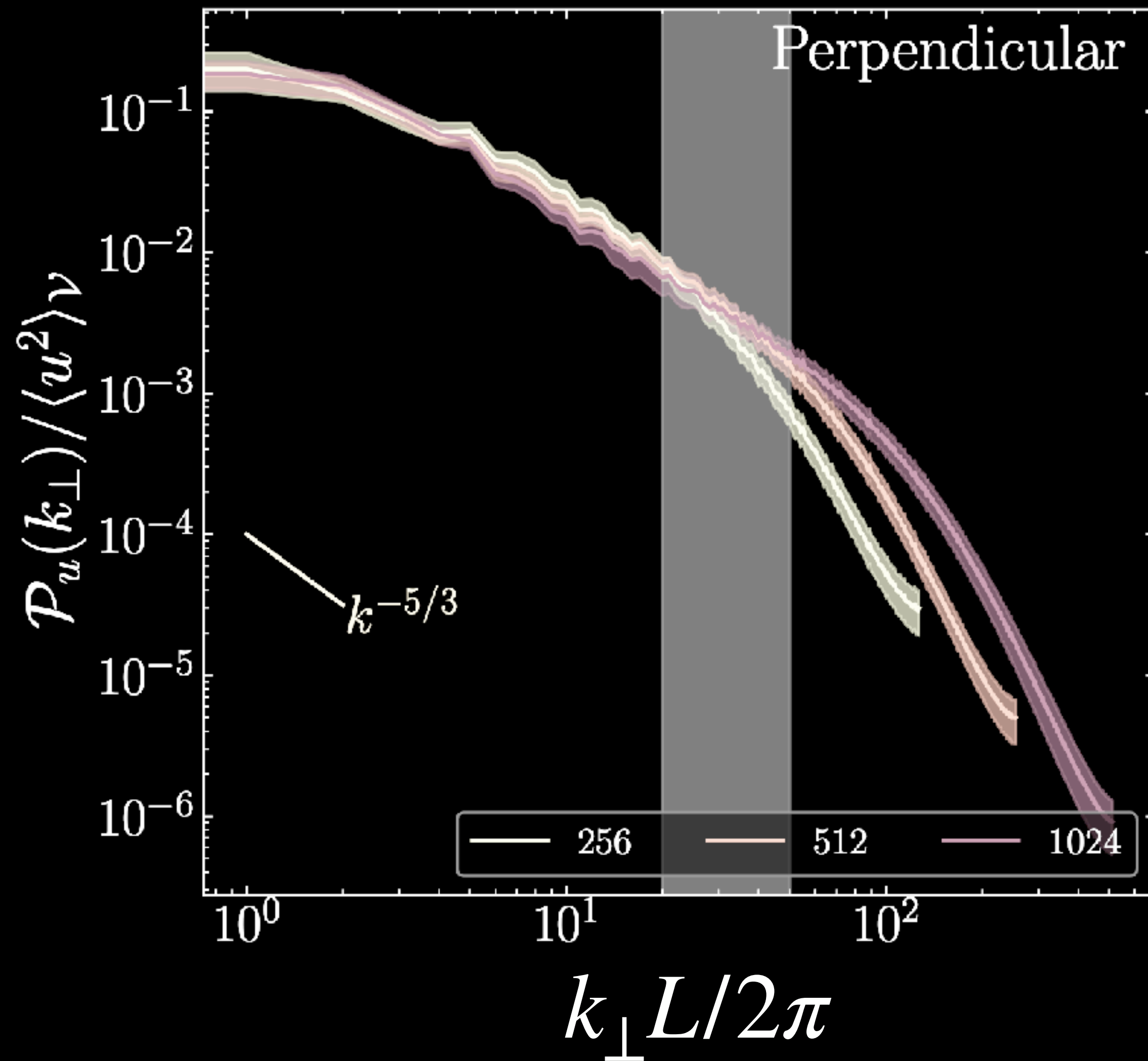


Fairly isotropic up and across $\nabla \phi$
Roughly Kolmogorov in total energy



Velocity Spectra

Compressible modes and solenoidal modes live very different lives!



Compressible modes and solenoidal modes live very different lives!
We should treat them differently!

$$\mathbf{k}' \xrightarrow{\mathbf{k}''} \mathbf{k}'''$$

cascade
transfers

$$\mathcal{T}_{cc}^c(k', k''') = - \int dV \mathbf{u}_c''' \otimes \mathbf{u}_c'' : \nabla \otimes \mathbf{u}_c'$$

$$\mathcal{T}_{ss}^s(k', k''') = - \int dV \mathbf{u}_s''' \otimes \mathbf{u}_s'' : \nabla \otimes \mathbf{u}_s'$$

•
•
•

$$\mathcal{T}_{cs}^s(k', k''') = - \int dV \mathbf{u}_s''' \otimes \mathbf{u}_s'' : \nabla \otimes \mathbf{u}_c'$$

interaction
transfers

Compressible modes and solenoidal modes live very different lives!
We should treat them differently!

$$\mathcal{T}_{cc}^c(k', k''') = - \int dV \mathbf{u}_c''' \otimes \mathbf{u}_c'' : \nabla \otimes \mathbf{u}_c'$$

$$\mathbf{k}' \xrightarrow{\mathbf{k}''} \mathbf{k}'''$$

cascade
transfers

$$\mathcal{T}_{ss}^s(k', k''') = - \int dV \mathbf{u}_s''' \otimes \mathbf{u}_s'' : \nabla \otimes \mathbf{u}_s'$$

•
•
•

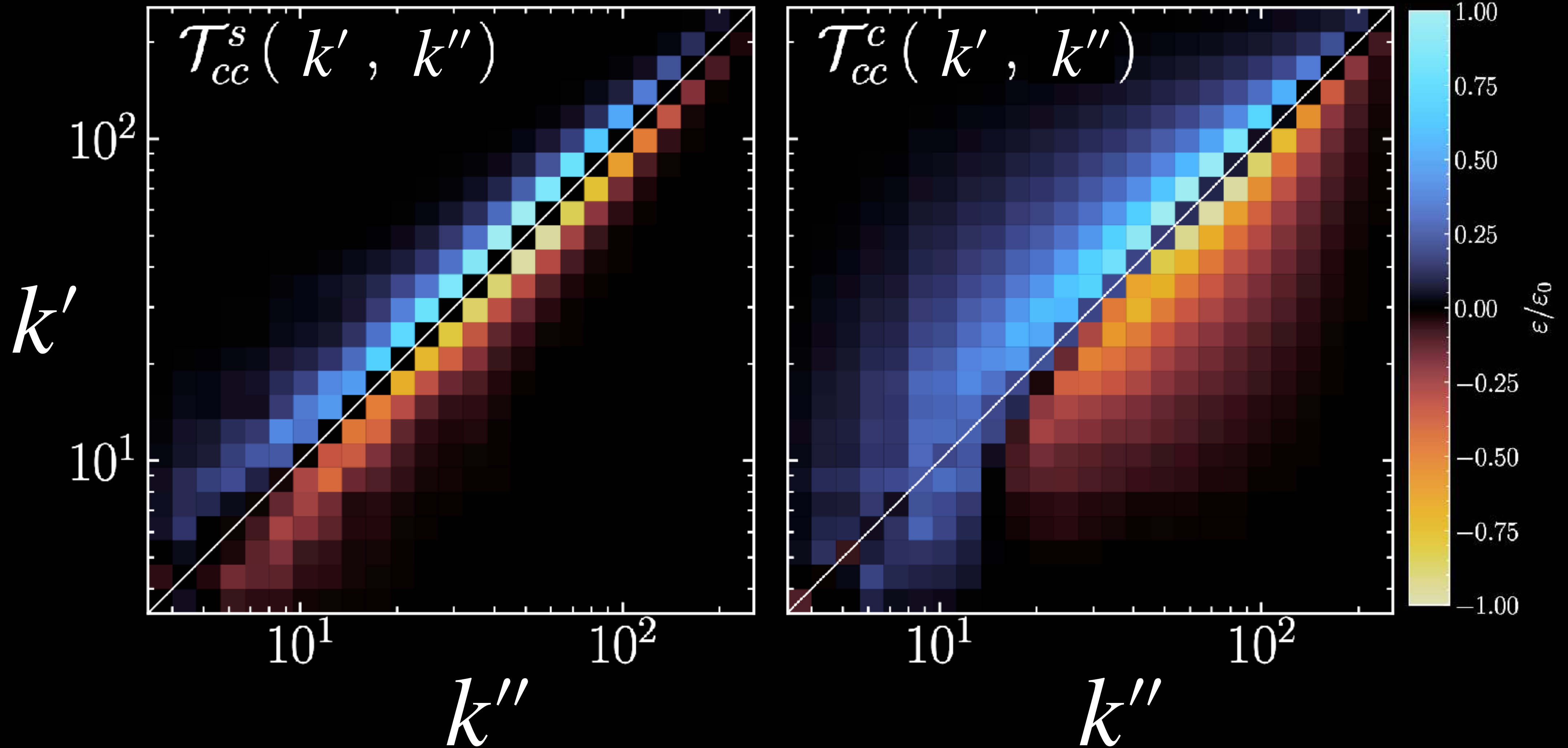
$$\mathcal{T}_{cs}^s(k', k''') = - \int dV \mathbf{u}_s''' \otimes \mathbf{u}_s'' : \nabla \otimes \mathbf{u}_c'$$

interaction
transfers

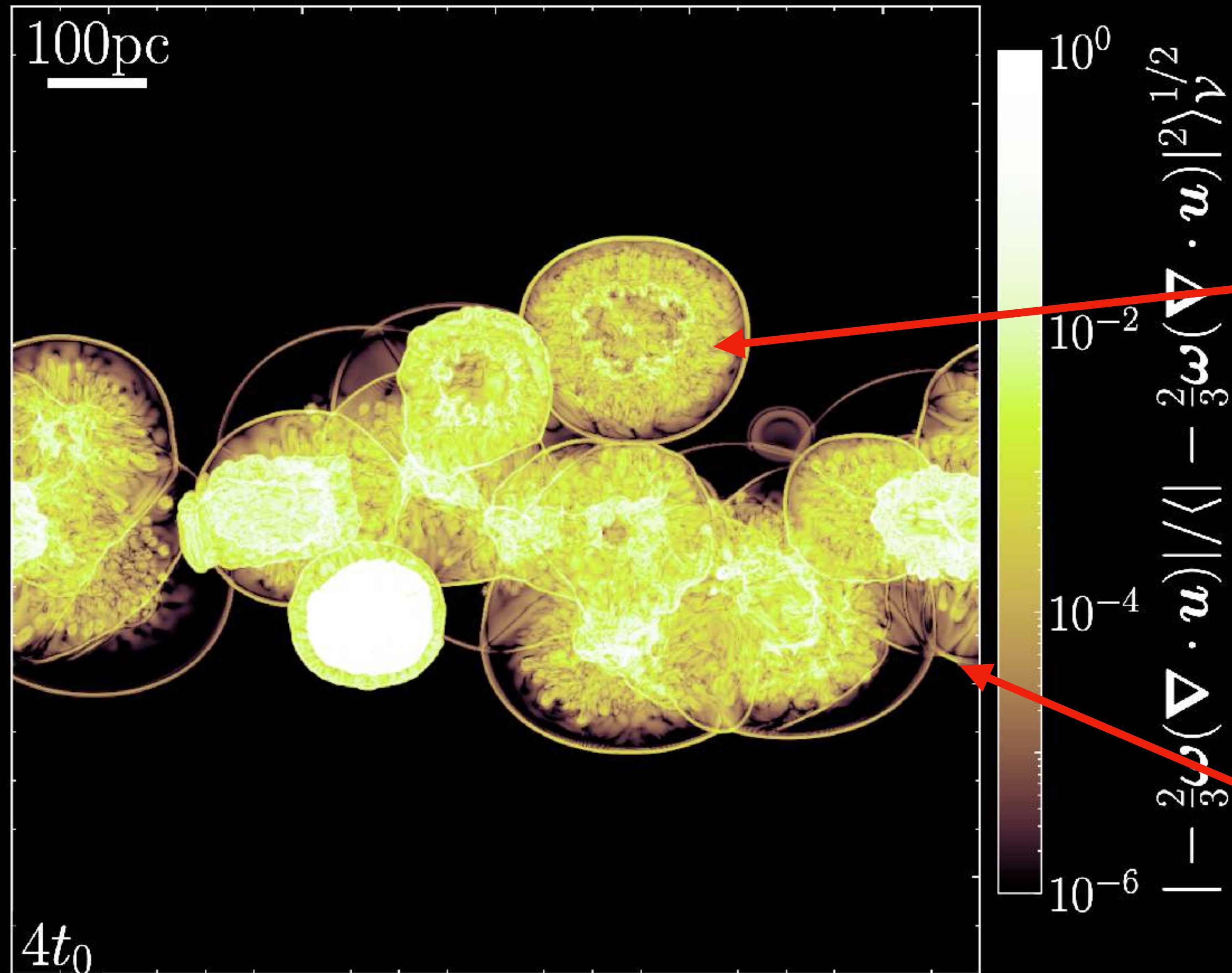
Cascade transfers: compressible modes

$$u'_c \xrightarrow{u''_s} u'''_c$$

$$u'_c \xrightarrow{u''_c} u'''_c$$



Cascade transfers: compressible modes



$$u'_c \xrightarrow{u''_s} u'''_c$$

Compressible modes
surfing solenoidal modes
down post-shock regions

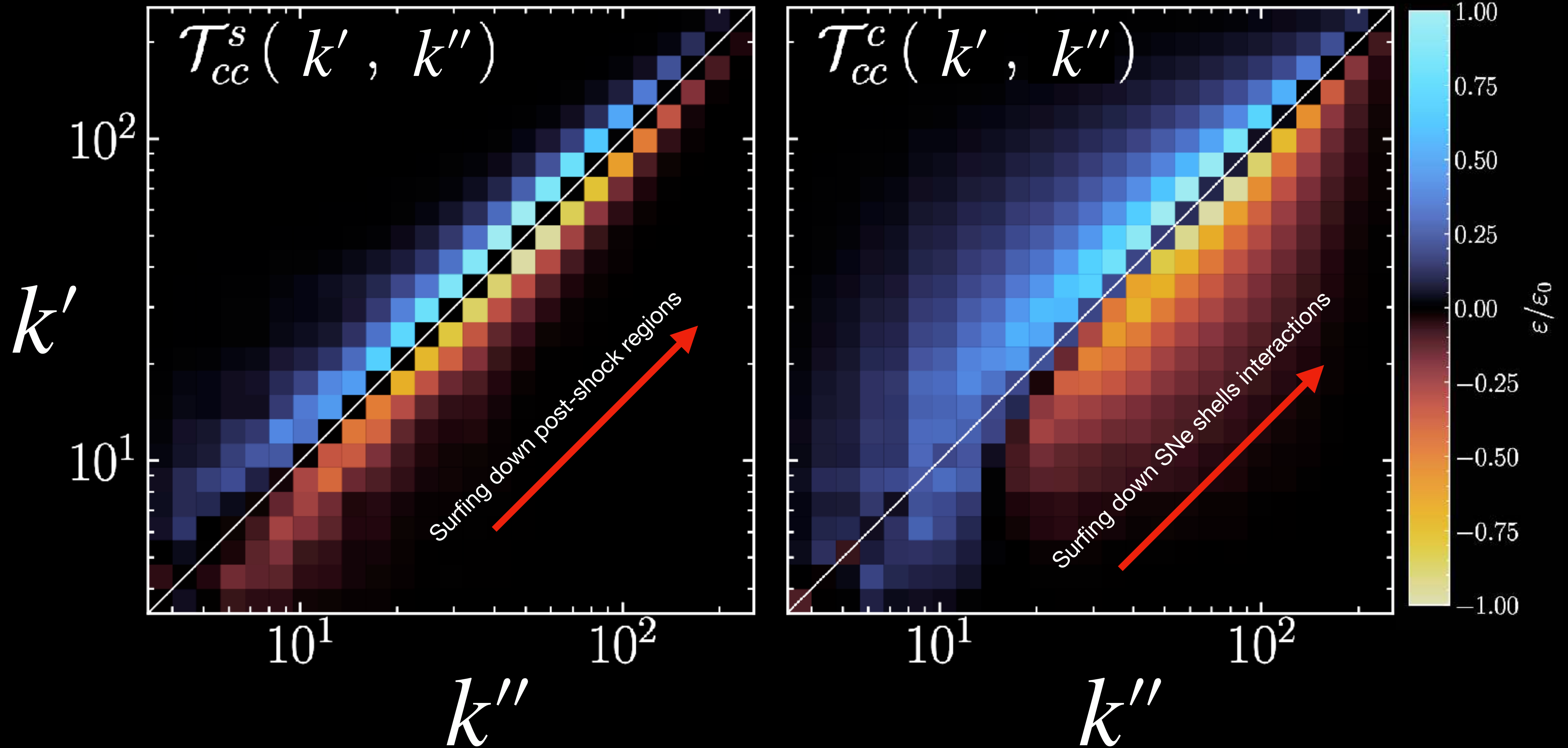
$$u'_c \xrightarrow{u''_c} u'''_c$$

Compressible modes
forming other compressible
modes through SNe shells

Cascade transfers: compressible modes

$$u'_c \xrightarrow{u''_s} u'''_c$$

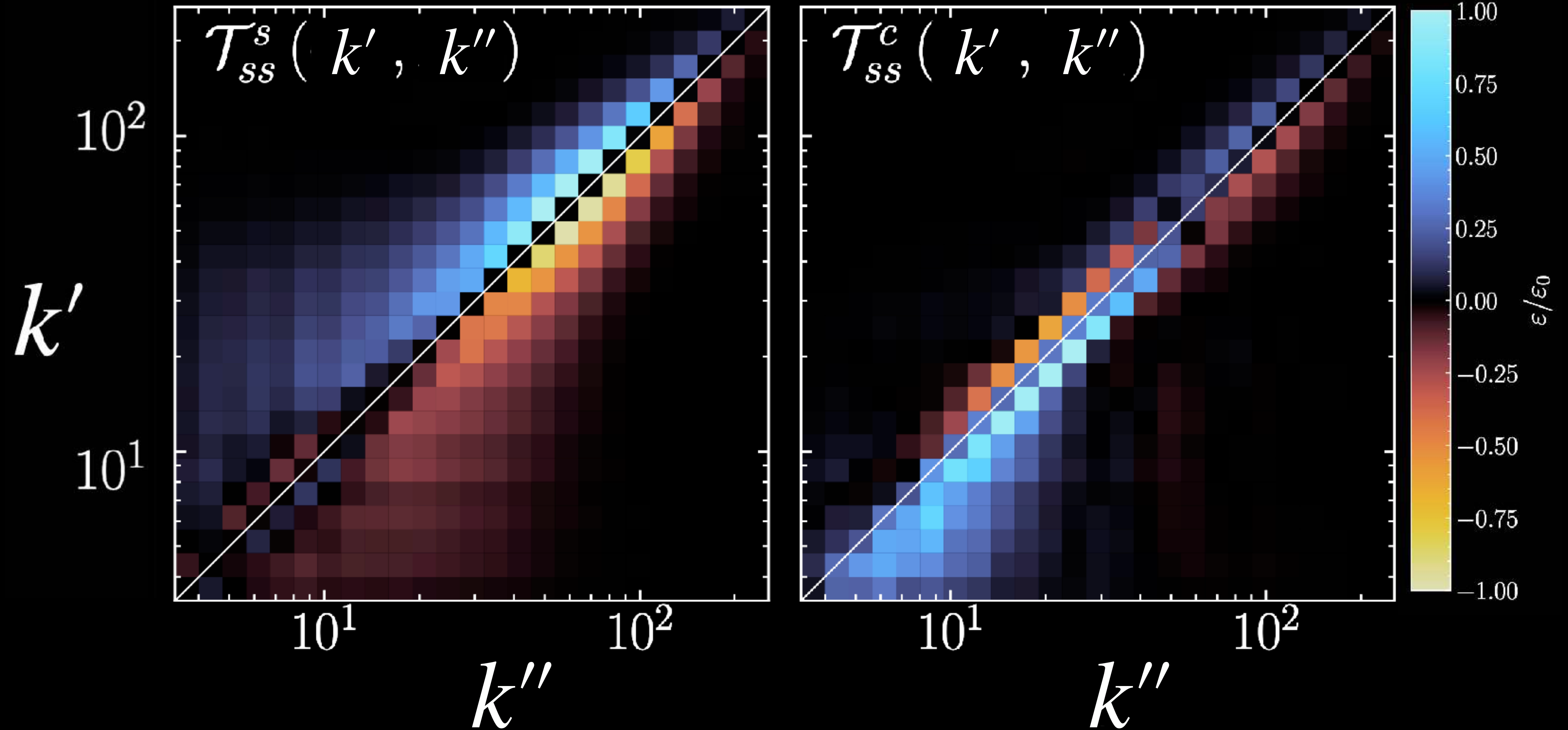
$$u'_c \xrightarrow{u''_c} u'''_c$$



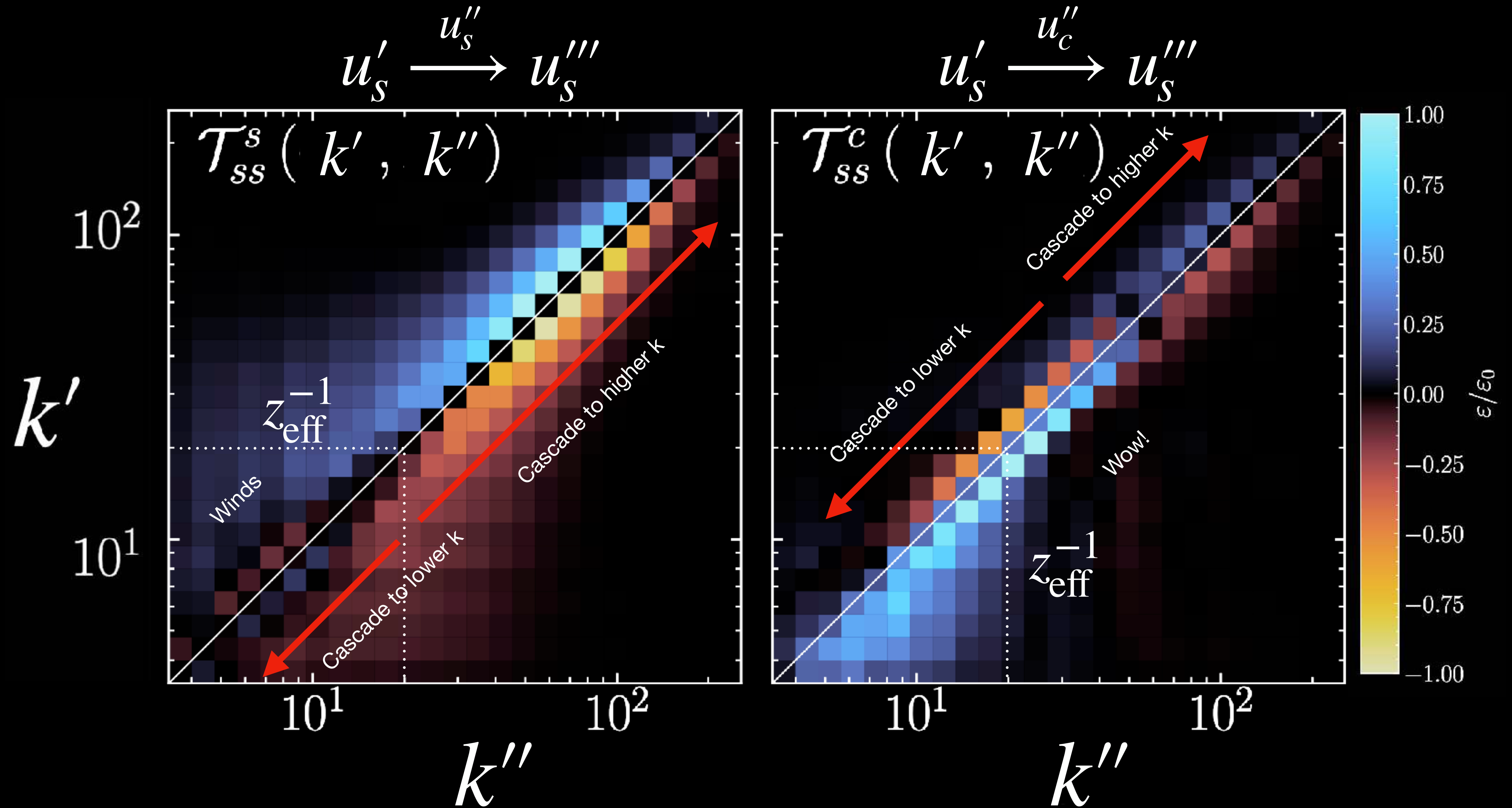
Cascade transfers: solenoidal modes

$$u'_s \xrightarrow{u''_s} u'''_s$$

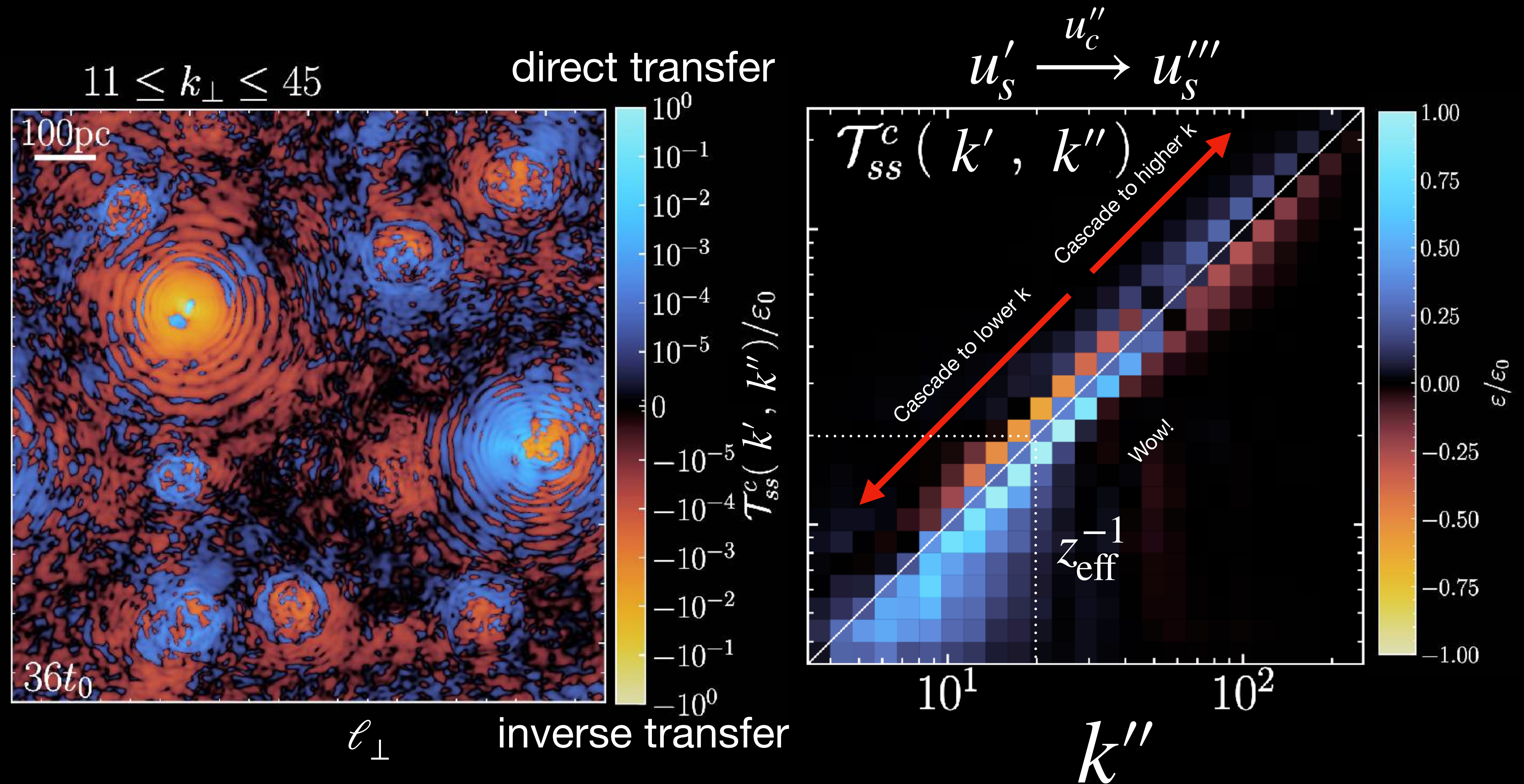
$$u'_s \xrightarrow{u''_c} u'''_s$$



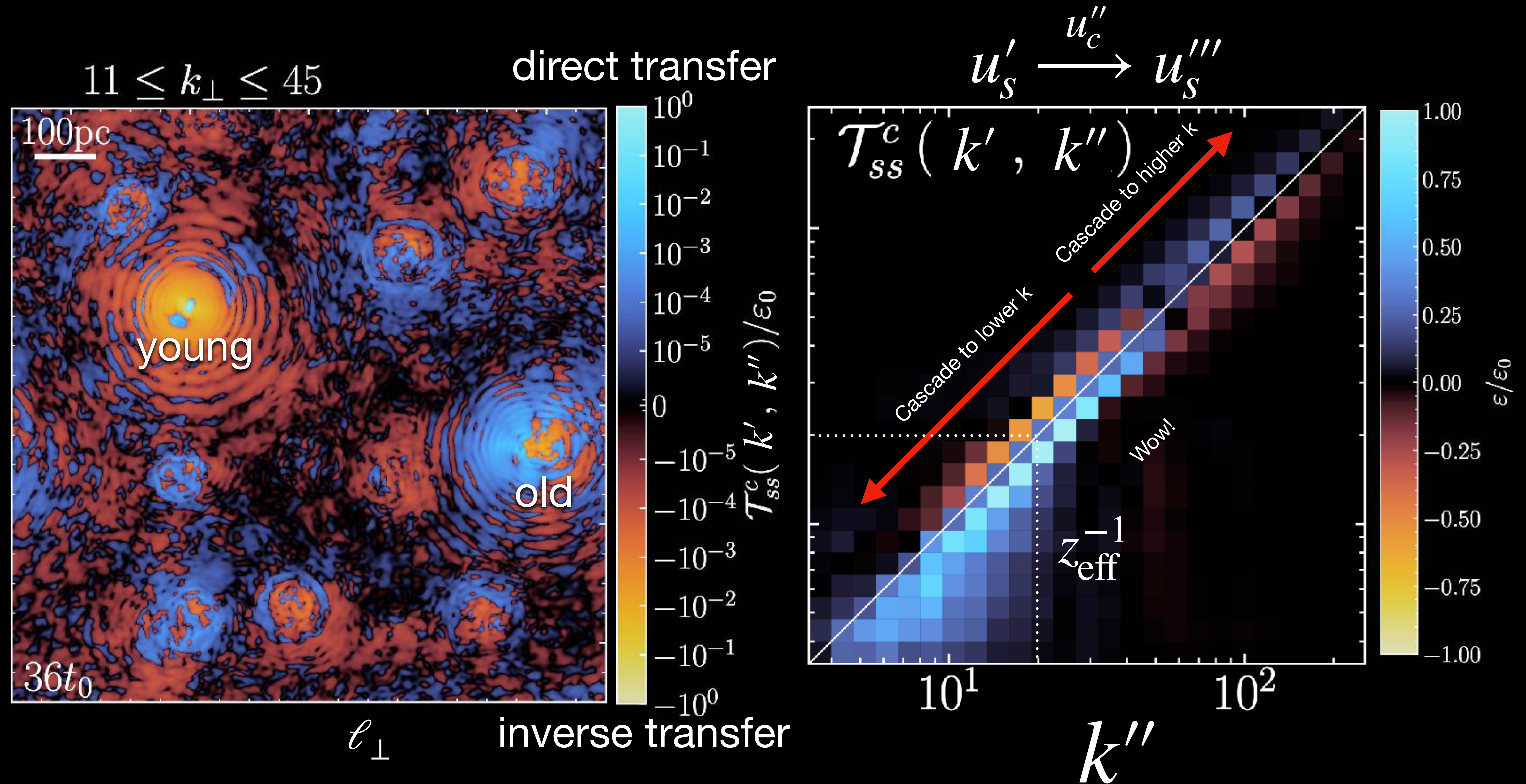
Cascade transfers: solenoidal modes



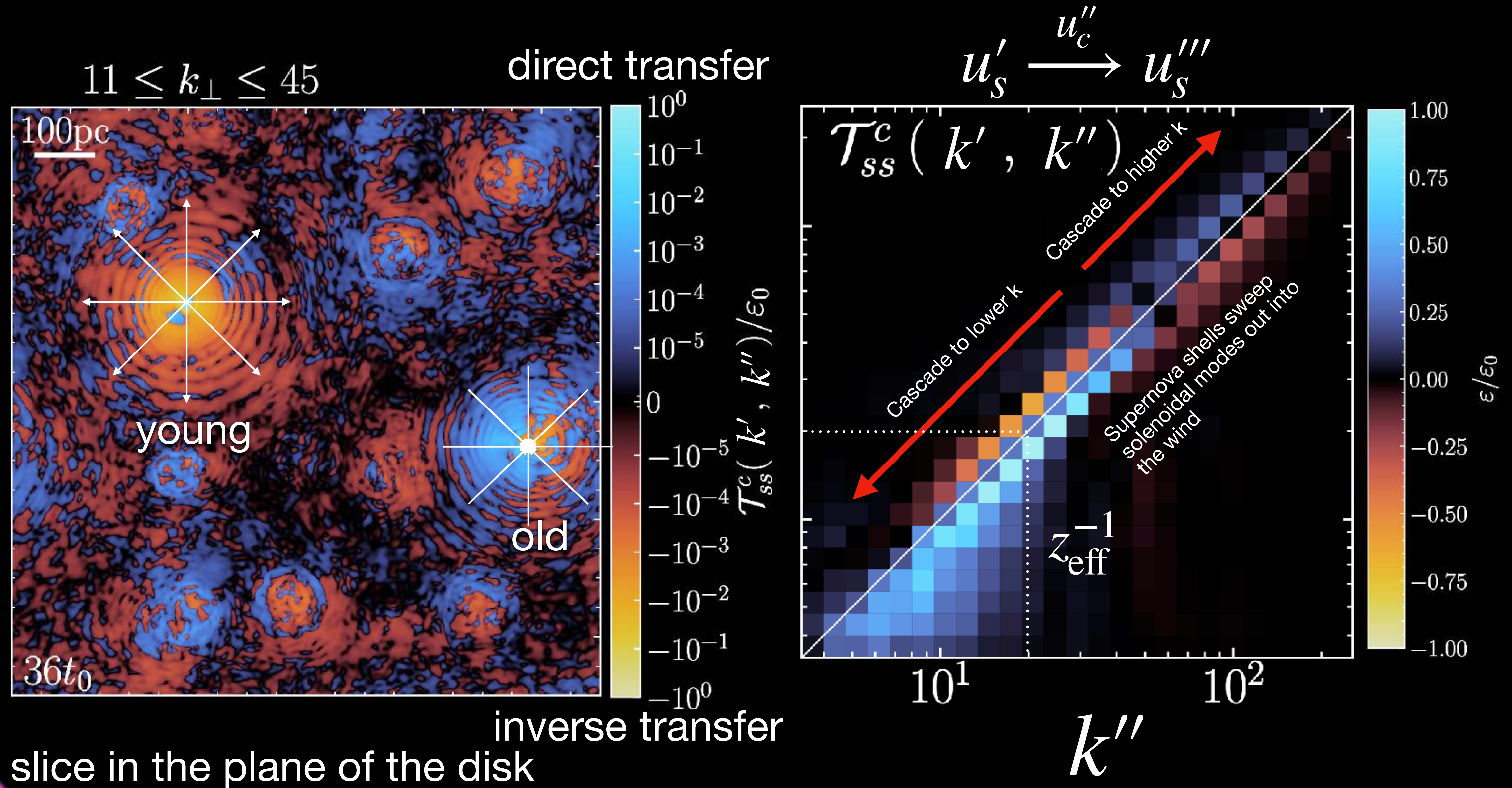
Cascade transfers: solenoidal modes



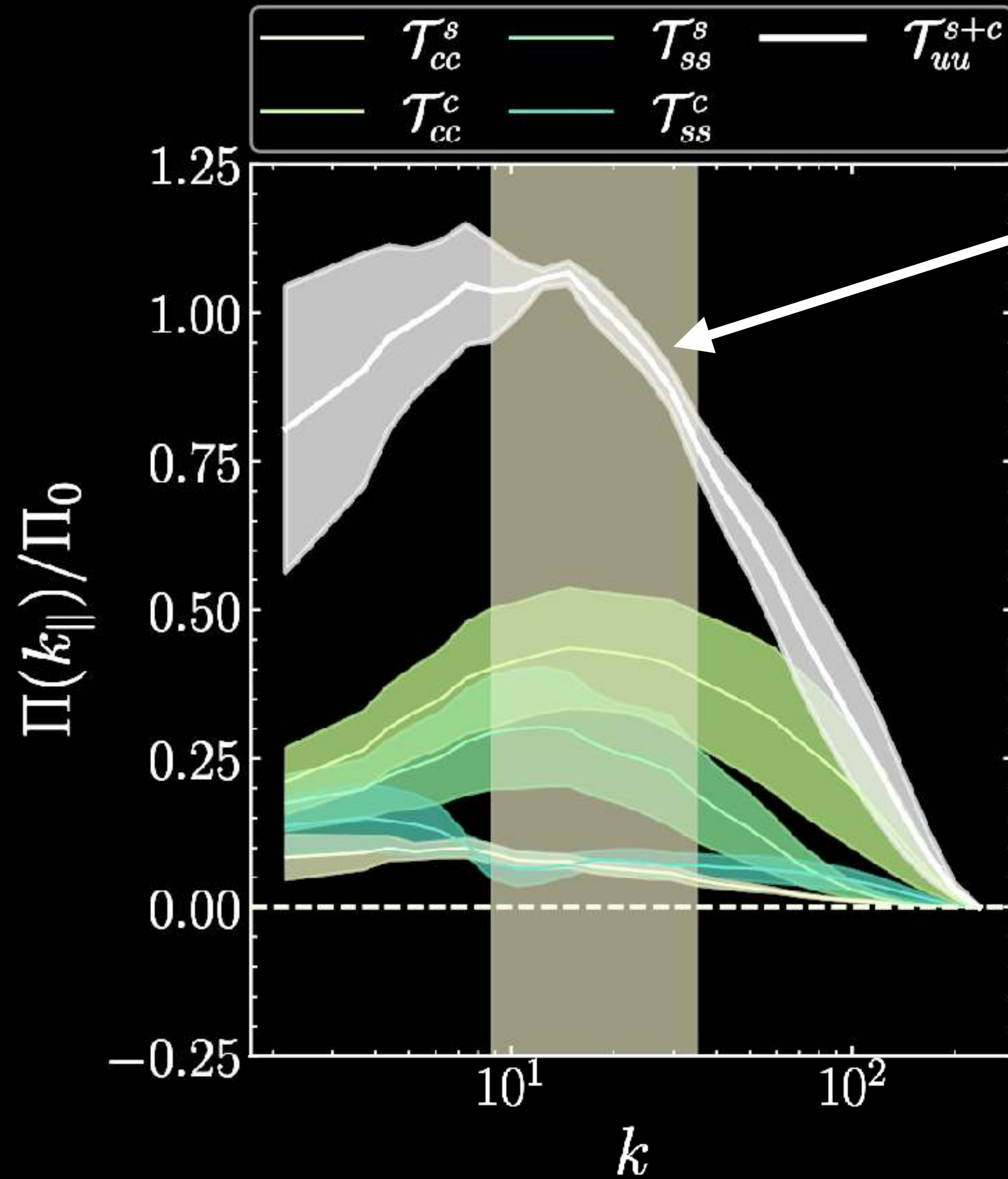
Cascade transfers: solenoidal modes — surfing on supernova shells



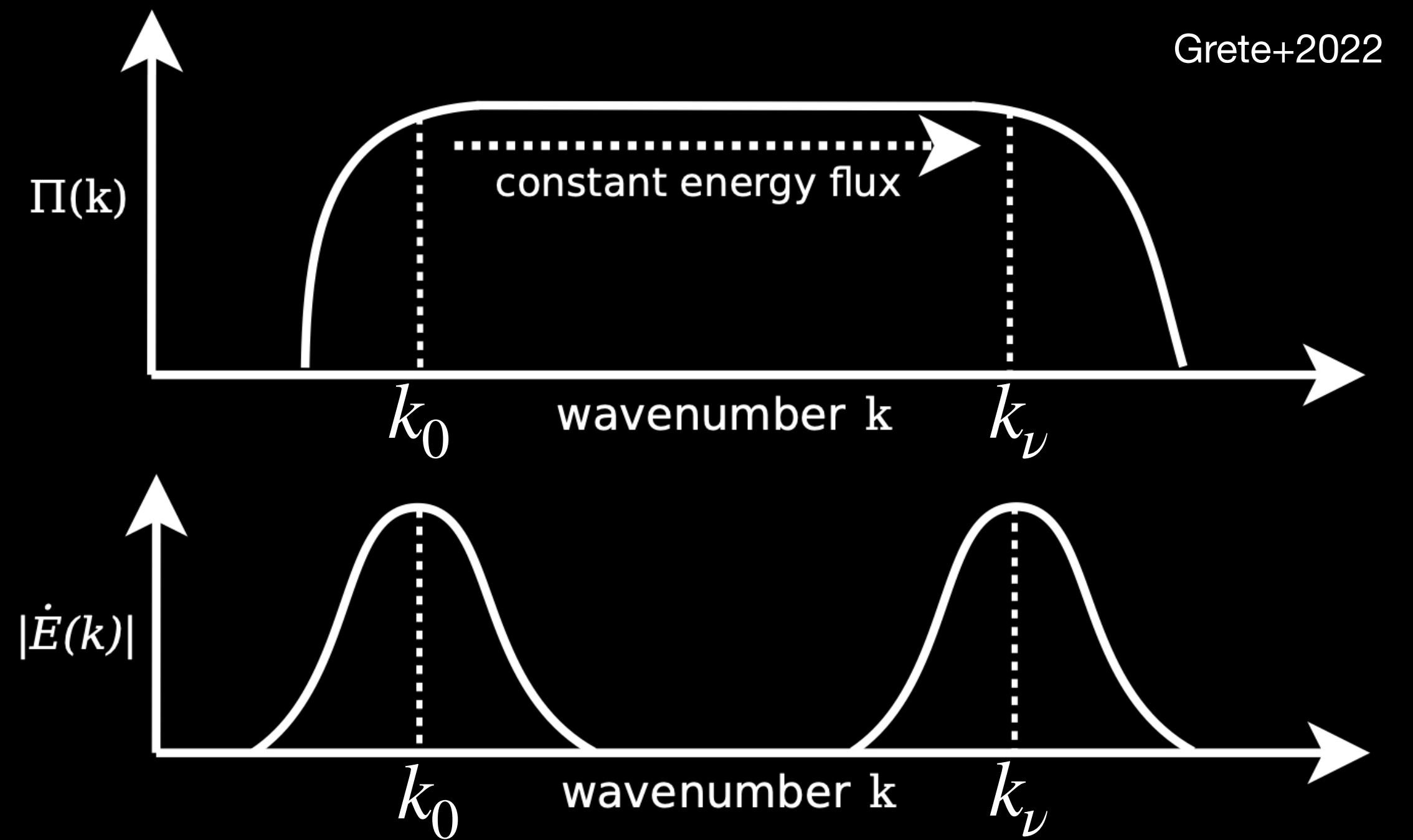
Cascade transfers: solenoidal modes — surfing on supernova shells



Cascade transfers: solenoidal modes — surfing on supernova shells



Not very constant at all!



Next steps with observers



james.beattie@princeton.edu



@astro_magnetism

- Mostly the density is accessed via observation... any hope to measure energy flux density? Yes!
- With sufficiently resolved column density (Francois, M-A ;)), one should be able to compute:

$$\mathcal{T}_{\rho\rho}(k', k''' | k'') = - \int dV \rho''' \mathbf{u}'' \cdot \nabla \otimes \rho' - \int dV \rho''' \rho' \nabla \cdot \mathbf{u}''$$

compression mediated cascade

advection mediated cascade

assuming isotropy. Never been done before...

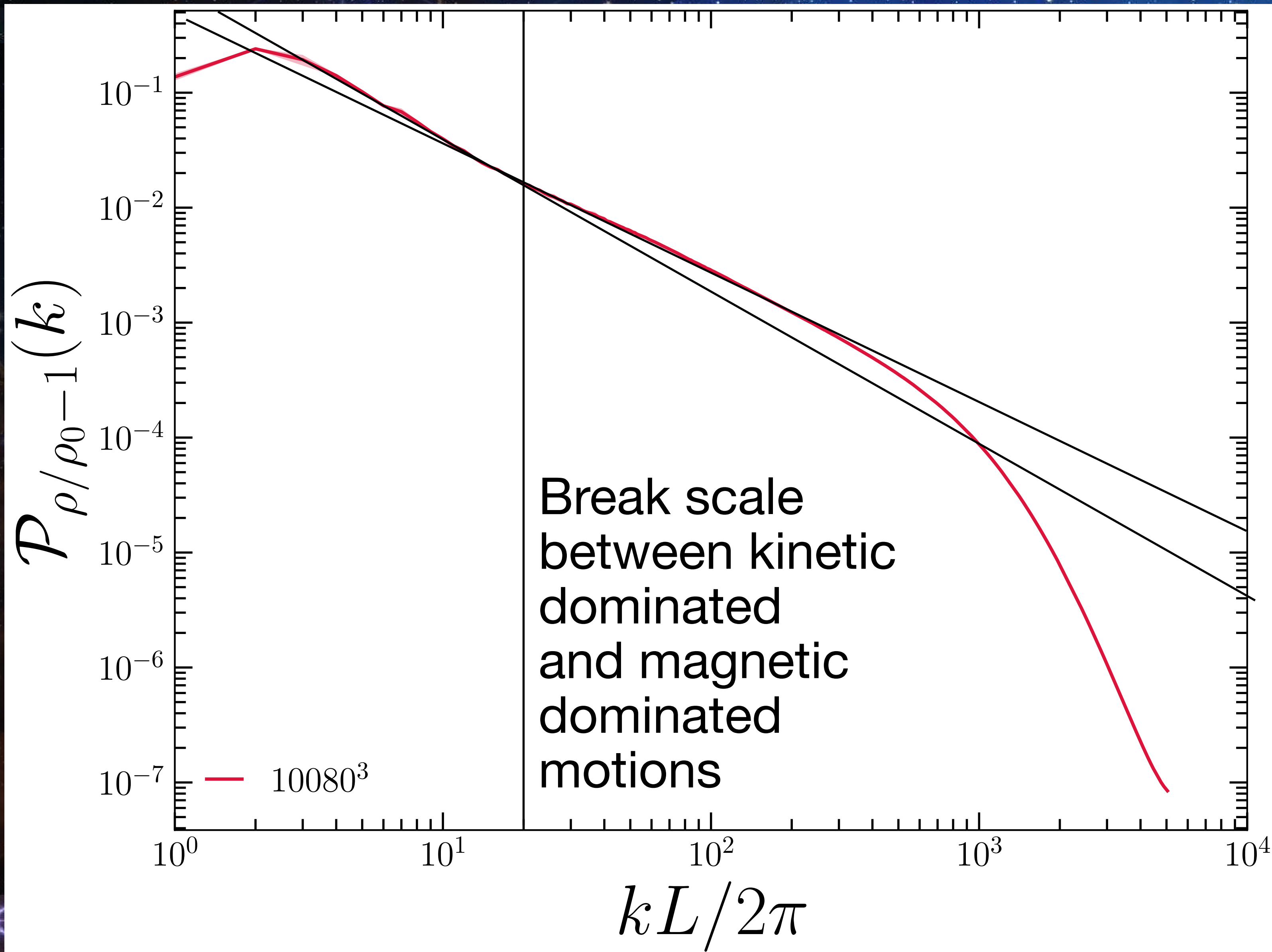
Next steps with observers



[james.beattie@princeton.](mailto:james.beattie@princeton.edu)



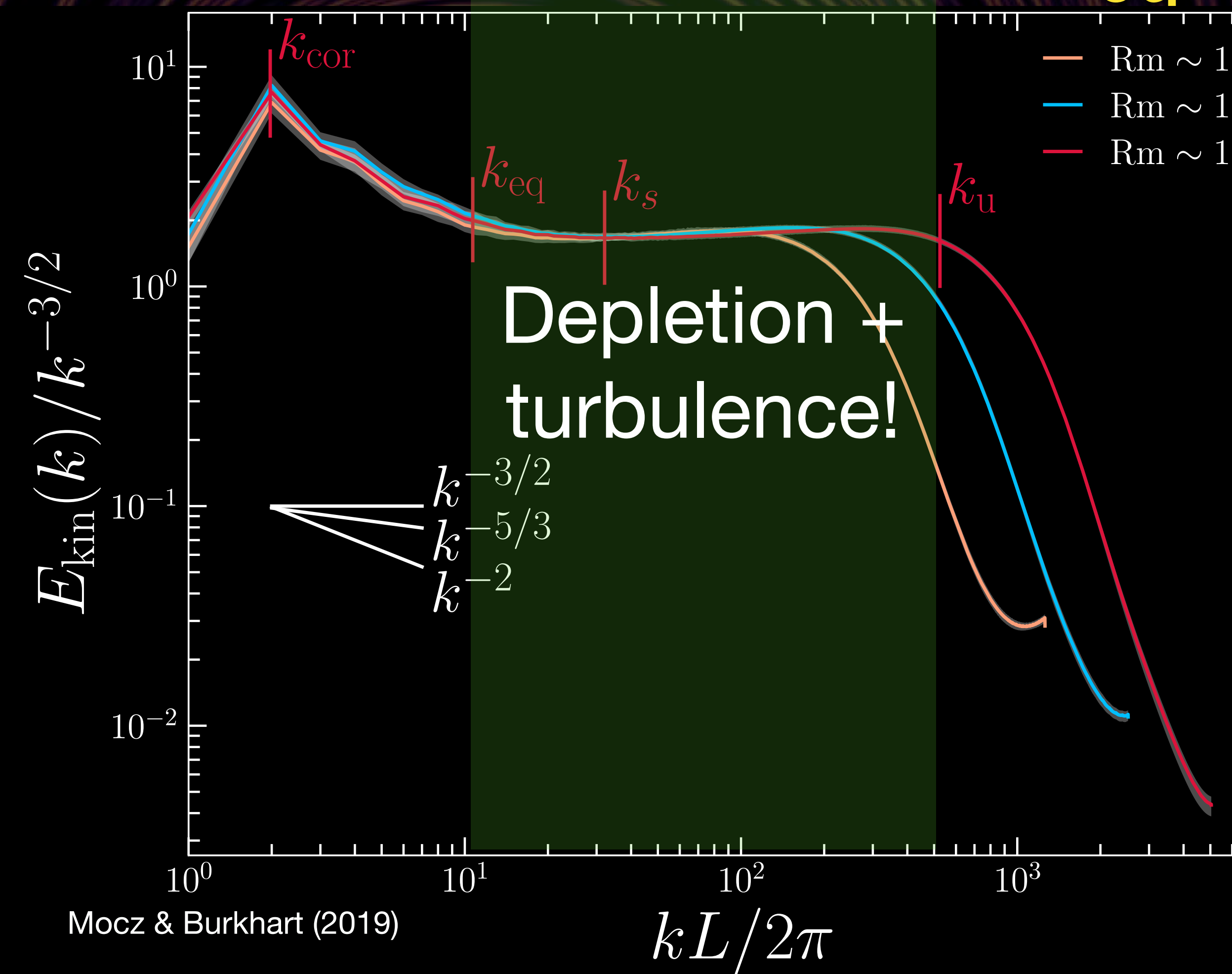
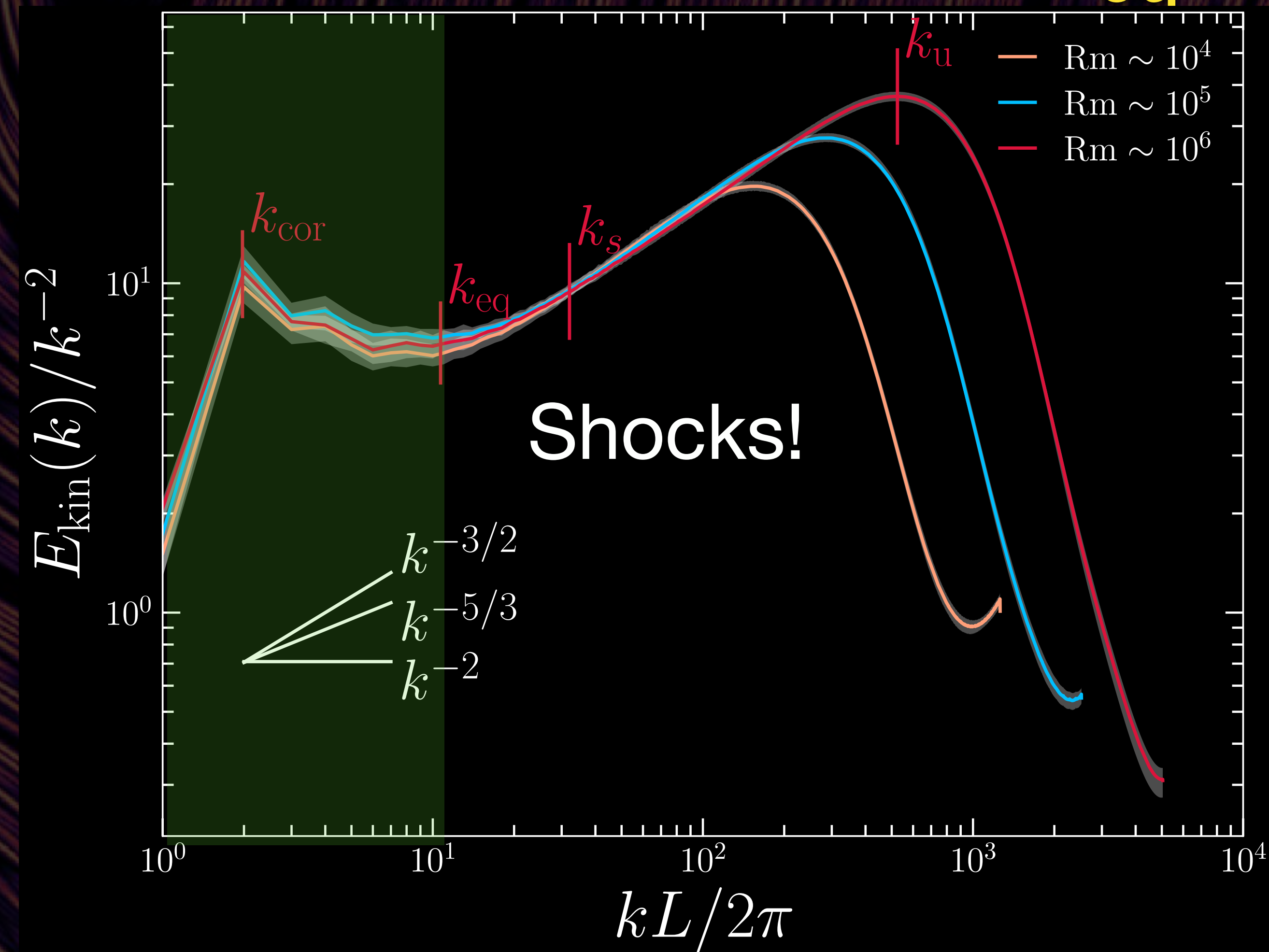
@astro_magnetism



Final thoughts (from Tuesday)...

$$\mathcal{E}(k) \sim k^{-2}, k \leq k_{\text{eq}}$$

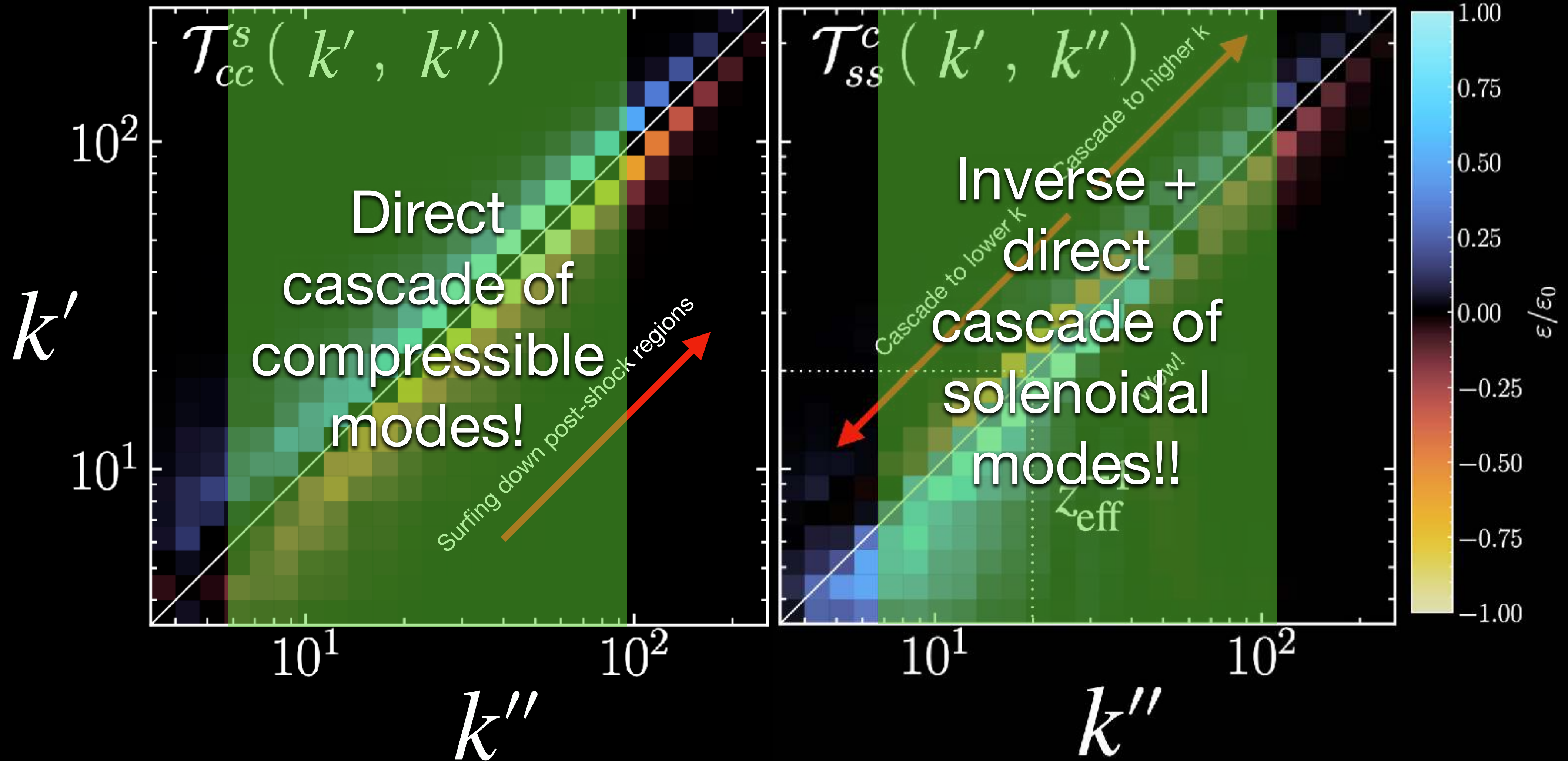
$$\mathcal{E}(k) \sim k^{-3/2}, k > k_{\text{eq}}$$



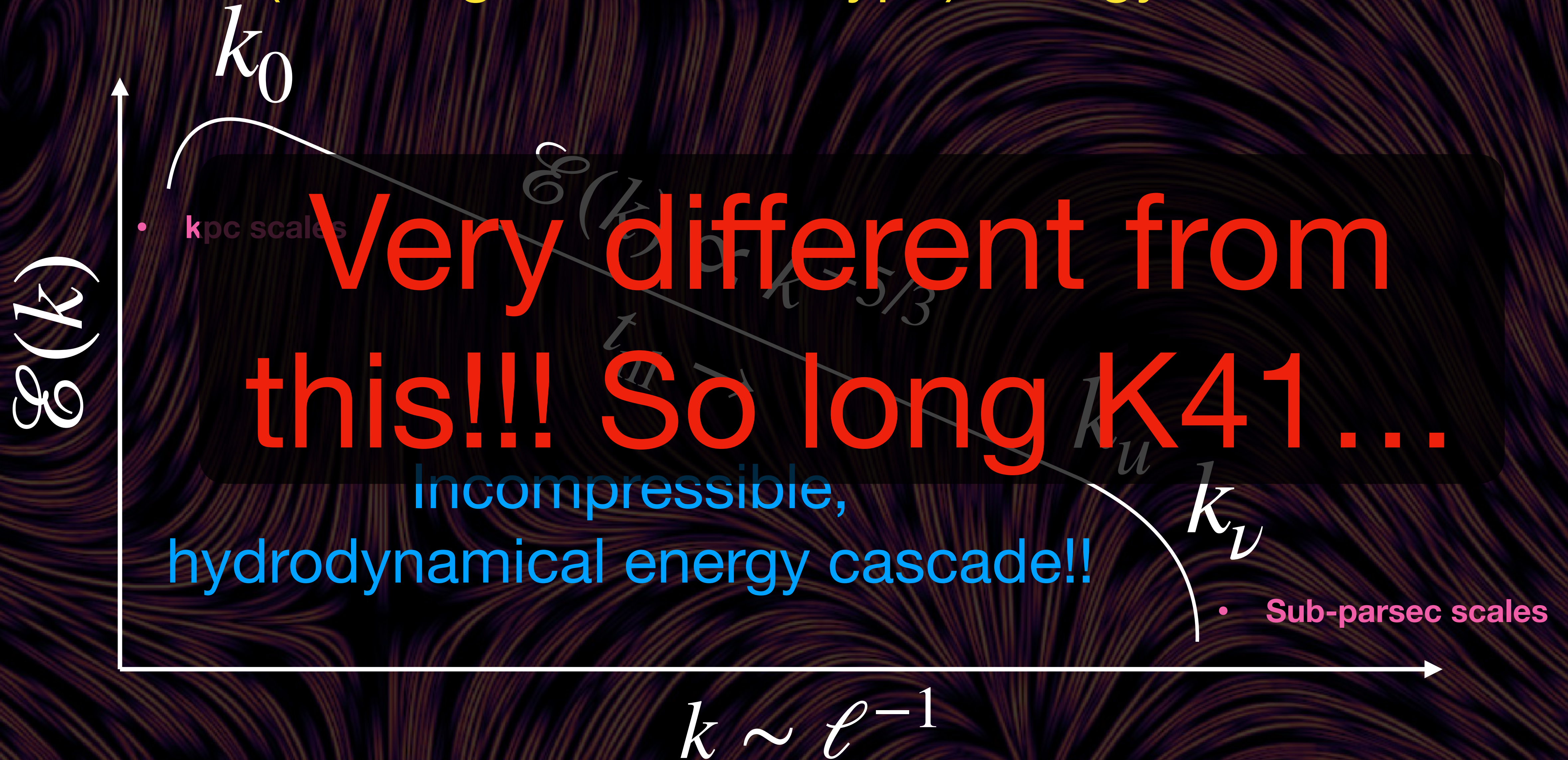
Final thoughts (from now)...

$$u'_c \xrightarrow{u''_s} u'''_c$$

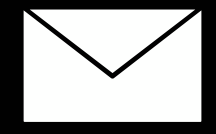
$$u'_s \xrightarrow{u''_c} u'''_s$$



The (Kolmogorov, 1941 -type) energy cascade



Thanks, questions?



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