

Can you create a large-scale vertical field from the KHI? Yup.

James R. Beattie^{1,2}, Amitava Bhattacharjee¹, Elias Most³, Bart Ripperda², Sasha Philippov⁴

¹Princeton University, ²Canadian Institute for Theoretical Astrophysics, ³Caltech, ⁴University of Maryland

Motivation:

Merging compact objects excite KH unstable shear layers in the plasma between the stars.

Due to a large separation in shear timescales t_Δ , and merger timescales, t_{merge} , shear layer dynamics can be realized. For example, in binary neutron star (BNS) mergers

$$t_\Delta = \frac{\delta_0}{\Delta U} \approx \frac{65 \text{ m}}{0.2c} = \mathcal{O}(10^{-6} \text{ s}) \approx 10^{-3} t_{\text{merge}}.$$

In $10^3 t_\Delta$, the shear layer can produce a strong magnetic through the magnetic dynamo. Tripathi et al. (2026) demonstrates the existence of a toroidal (in-plane) dynamo. Gutiérrez et al. (2026) shows a strong fluctuation dynamo. But most important for multimessenger signals of the events is the BZ mechanism, which requires a vertical field dynamo. Here we show, for the first time, a vertical field dynamo that is active in the first $500 t_\Delta$ in the shear layer.

Numerical Setup:

3D, visco-resistive, compressible MHD equation, with KHI ICs:

$$\frac{u_y(x, t_0)}{U_y} = \Delta(x) - 1, \quad \Delta U = 2U_y, \quad \text{shear profile}$$

$$\frac{\rho(x, t_0)}{\rho_0} = 1 + \frac{\delta\rho}{2\rho_0} \Delta(x), \quad \text{stratification}$$

$$\Delta(x) = \tanh\left(\frac{x-x_1}{\delta_0}\right) - \tanh\left(\frac{x-x_2}{\delta_0}\right),$$

$$\frac{u_x(x, y, t_0)}{U_x} = \frac{1}{\sqrt{n_{\text{max}} + 1}} \sum_n^{n_{\text{max}}} \sin\left(\frac{2n\pi}{L}y + \phi_n\right) \times \text{perturbation} \\ \left(\exp\left\{-\frac{(x-x_1)^2}{\sigma^2}\right\} + \exp\left\{-\frac{(x-x_2)^2}{\sigma^2}\right\} \right),$$

with initial random magnetic field that satisfies:

$$\tilde{B}_i(\mathbf{k}) = \tilde{A}(k)(\delta_{ij} - k_i k_j / k^2) \tilde{g}_j(\mathbf{k}) \iff \mathbf{k} \cdot \tilde{\mathbf{B}}(\mathbf{k}) = 0, \quad k \neq 0,$$

$$\langle \tilde{g}_i(\mathbf{k}) \rangle = 0, \quad \langle \tilde{g}_i(\mathbf{k}) \tilde{g}_j^\dagger(\mathbf{k}') \rangle = \delta_{ij} \delta_{\mathbf{k}\mathbf{k}'}. \quad \text{magnetic field breaks } z \text{ symmetry}$$

discretized on uniform grids: $288^3 \leq N_{\text{grid}}^3 \leq 4,608^3$

Dynamo Results:

Integral energy statistics are dominated by large-scale, diffusing shear layer:

$$\frac{d \ln E_{\text{mag}}}{dt} = 2 \hat{B}_i \hat{B}_j \partial_j u_i \sim \frac{2U_y}{\delta(t)} \sim \frac{2U_y}{\sqrt{\delta_0^2 + C_\nu \nu t}},$$

$$\text{and integrating} \\ \frac{E_{\text{mag}}(t)}{E_{\text{mag}}(0)} = \exp\left(\frac{2C_d \text{Re}_\Delta}{C_\nu} \left[\sqrt{1 + \frac{C_\nu}{\text{Re}_\Delta} \frac{t}{t_\Delta}} - 1 \right]\right),$$

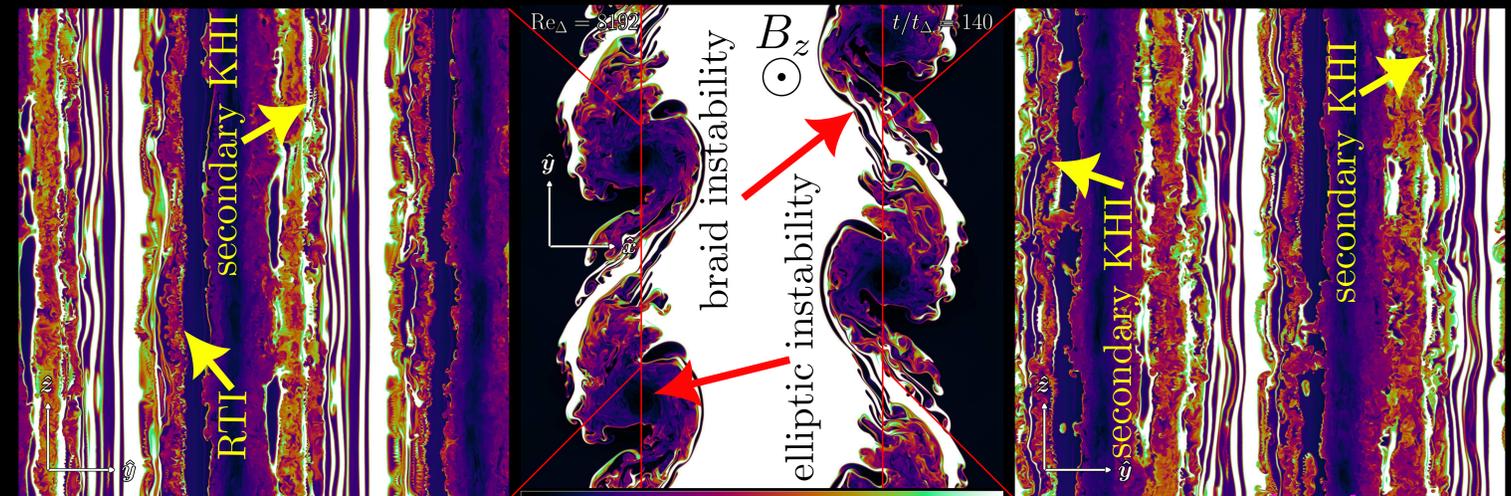
where $C_d = \mathcal{O}(10^{-1})$ is the dynamo efficiency.

For $t/t_\Delta \ll \text{Re}_\Delta / C_\nu$,

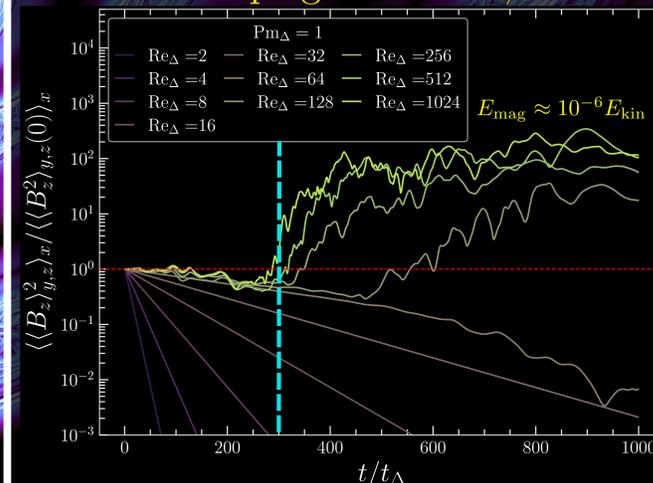
$$\frac{E_{\text{mag}}}{E_{\text{mag}}(0)} \approx \exp\left(C_d \frac{t}{t_\Delta}\right),$$

predicts asymptotic $\text{Re} \rightarrow \infty$ growth rate, $\gamma \approx C_d / t_\Delta$.

KHI goes nonlinear on $t \gtrsim 1t_\Delta$, creating roll-ups. Roll-ups go 3D unstable via a zoo of instabilities:



As the roll-ups go unstable, a vertical field dynamo ignites



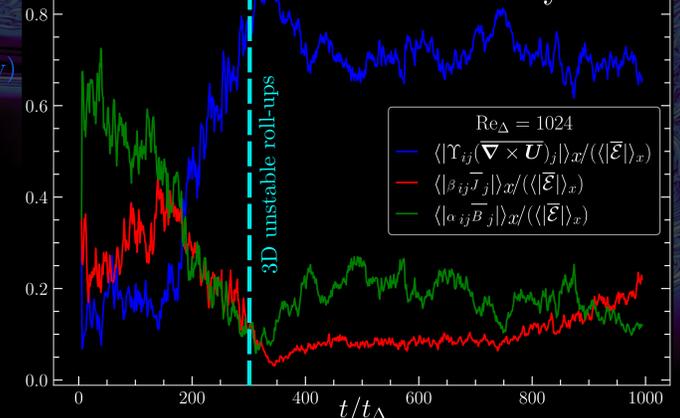
Upsilon term
(inhomogenous velocity)
Tripathi et al. (2026)

70% flux

Beta term
70% flux

Alpha term
70% flux

EMF terms for vertical field dynamo



The vertical field dynamo has a mix of EMF contributions, but is dominated by velocity inhomogeneities.